

MAE263B: Dynamics of Robotic Systems Discussion Section – Week6

: Jacobian (SCARA)

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Homework

Homework - Review

- On CCLE, you can check your score and download an annotated pdf version for your homework.
- All reasons for deduction are described in the pdf file.

Solution will be uploaded on CCLE too

- The solution for homework2 on the bionics lab course website is wrong.
- The solution for homework3 is not uploaded on the course website.
- \rightarrow From this week, solution files will be uploaded on CCLE too.

Contents

□ Jacobian with SCARA example

- Force/Torque propagation
- Review
 - Velocity propagation
 - Direct differentiation



Kinematics Relations - Joint & Cartesian Spaces

- The location of the robot end-effector may be specified using one of the following descriptions:
- Cartesian Space / Operational Space









Kinematics Relations - Forward & Inverse

• The robot kinematic equations relate the two description of the robot tip location







• The velocity relationship: The relationship between the joint angle rates (\dot{X}) and the linear and angular velocities of the end effector $(\dot{\theta})$.

 $\dot{\mathbf{v}} = \mathbf{i}\dot{\mathbf{o}}$



This expression can be expanded to: $\underline{\dot{x}} = J(\underline{\theta})\underline{\dot{\theta}}$ •



- ٠
 - is a 6x1 vector of the end effector linear and angular velocities X
 - is a 6xN Jacobian matrix $-J(\theta)$
 - is a Nx1 vector of the manipulator joint velocities θ_{N}
 - is the number of joints N

The velocity relationship

: The relationship between the joint angle rates ($\underline{\dot{\theta}}_N$) and the translation and rotation velocities of the end effector ($\underline{\dot{x}}$).

$$\underline{\dot{x}} = J(\underline{\theta})\underline{\dot{\theta}}$$

 The relationship between the robot joint torques (<u>τ</u>) and the forces and moments (<u>F</u>) at the robot end effector (Static Conditions).

$$\underline{\tau} = J(\underline{\theta})^T \underline{F}$$







• This expression $(\underline{\tau} = J(\underline{\theta})^T \underline{F})$ can be expanded to:



- Where:
 - $\frac{\tau}{2}$ is a Nx1 vector of the robot joint torques
 - $-J(\underline{\theta})^{T}$ is a Nx6 Transposed Jacobian matrix
 - <u>F</u> is a 6x1 vector of the forces and moments at the robot end effector
 - N is the number of joints





Jacobian Matrix - Derivation Methods







Jacobian – static force & torque

In addition to the velocity relationship, we are also interested in developing a relationship between the robot joint torques (<u><u>r</u></u>) and the forces and moments (<u>F</u>) at the robot end effector (Static Conditions). This relationship is given by:

$$\begin{bmatrix} \tau_1 \\ \tau_2 \\ \dots \\ \tau_N \end{bmatrix} = \begin{bmatrix} J(\underline{\theta}) \\ M_x \\ M_y \\ M_z \end{bmatrix}$$

$$\underline{\tau} = J(\underline{\theta})^{\underline{T}}\underline{F}$$







Jacobian – static force & torque

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FIGURE 6.1. Fanue LR Mate robot. (Courtesy of Fanue Robotics North America, Inc., Rochester Hills, Michigan.)

• The force / moment equilibrium equations $\sum F = 0$ $\sum M = 0$

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- Typically, the robot is **pushing on something** in the environment with the end-effector or is perhaps **supporting a load** at the hand.
- We wish to solve for the joint torques that must be acting to keep the system in static equilibrium.
- The force / moment equilibrium equations

$$\mathbf{F} \quad \sum \mathbf{M} = \mathbf{0}$$





FIGURE 6.1. Fanue LR Mate robot. (Courtesy of Fanue Robotics North America, Inc., Rochester Hills, Michigan.)



- Static force analysis is of practical importance in determining the quality of force transmission through the various joints of a mechanism.
- It serves as a basis for sizing the links and bearings of a robot manipulator and for selecting appropriate actuators.
- The result also be used for force control of robot manipulator.





FIGURE 6.1. Fanue LR Mate robot. (Courtesy of Fanue Robotics North America, Inc., Rochester Hills, Michigan.)



Problem

Given: Typically the robot's end effector is applying forces and torques on an object in the environment or carrying an object (gravitational load).

Compute: The joint torques must be acting to keep the system in static equilibrium.







Solution

Jacobian - Mapping from the joint force/torques - $\underline{\tau}$ to forces/torque in the Cartesian space applied on the end effector) - \underline{F} .

$$\underline{\tau} = J(\underline{\theta})^T \underline{F}$$

Free Body Diagram - The chain like nature of a manipulator leads to decompose the chain into individual links and calculate how forces and moments propagate from one link to the next.







Static Analysis Protocol - Free Body Diagram 1/

Step 1

Lock all the joints - Converting the manipulator (mechanism) to a structure

Step 2

Consider each link in the structure as a free body and write the force / moment equilibrium equations

 $\begin{array}{ll} \mbox{(3 Eqs.)} & \sum F = 0 \\ \mbox{(3 Eqs.)} & \sum M = 0 \end{array}$

Step 3

Solve the equations - 6 Eq. for each link.

Apply backward solution starting from the last link (end effector) and end up at the first link (base)







Static Analysis Protocol - Free Body Diagram 2/

- Special Symbols are defined for the force and torque exerted by the neighbor link
 - f_i Force exerted on link i by link i-1
 - n_i Torque exerted on link i by link i-1

Reference coordinate system {B}



Exerted on link A by link A-1

 $\{i+1\}$ $\{i\}$ iP_{i+1} iP_{i+1} link i + 1 link i

For easy solution superscript index
 (B) should the same as the subscript
 (A)



Static Analysis Protocol - Free Body Diagram 3/

• For serial manipulator in static equilibrium (joints locked), the sum the forces and torques acting on Link i in the link frame $\{i\}$ are equal to zero.



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Static Analysis Protocol - Free Body Diagram 4/

• For serial manipulator in static equilibrium (joints locked), the sum the forces and torques acting on Link i in the link frame $\{i\}$ are equal to zero.





• Changing the reference frame such that each force (and torque) is expressed upon their link's frame, we find the static force (and torque) propagation from link i+1 to link i

$$\underbrace{\stackrel{i}{f_{i}} = \stackrel{i}{f_{i+1}}}_{i} = \stackrel{i}{i} \stackrel{i}{R} \stackrel{i+1}{f_{i+1}} f_{i+1}$$

$$\underbrace{\stackrel{i}{n_{i}} = \stackrel{i}{n_{i+1}} + \stackrel{i}{P_{i+1}} \times \stackrel{i}{f_{i+1}} = \stackrel{i}{i} \stackrel{i}{R} \stackrel{i+1}{n_{i+1}} + \stackrel{i}{P_{i+1}} \times \stackrel{i}{f_{i}} f_{i}$$

$${}^{i}f_{i} = {}^{i}_{i+1}R {}^{i+1}f_{i+1}$$

 ${}^{i}n_{i} = {}^{i}_{i+1}R {}^{i+1}n_{i+1} + {}^{i}P_{i+1} \times {}^{i}f_{i}$

- These equations provide the static force (and torque) propagation from link to link.
- They allow us to start with the force and torque applied at the end effector and calculate the force and torque at each joint all the way back to the robot base frame.





Static Analysis Protocol - Free Body Diagram 6/

• **Question:** What torques are needed at the joints in order to balance the reaction moments acting on the link (**Revolute Joint**).







- Question: What torques are needed at the joints in order to balance the reaction moments acting on the link (Revolute Joint).
- **Answer:** All the components of the moment vectors are resisted by the structure of mechanism itself, except for the torque about the joint axis (**Revolute Joint**).







Static Analysis Protocol - Free Body Diagram 8/

• **Question:** What forces are needed at the joints in order to balance the reaction forces acting on the link (**Prismatic Joint**).







Static Analysis Protocol - Free Body Diagram 9/

- **Question:** What forces are needed at the joints in order to balance the reaction forces acting on the link (**Prismatic Joint**).
- **Answer:** All the components of the force vectors are resisted by the structure of mechanism itself, except for the force along the joint (**Prismatic joint**).



$$f_{i} = {}^{i}f_{i}^{T}\hat{Z}_{i}$$
$$= [{}^{i}f_{ix} {}^{i}f_{iy} {}^{i}f_{iz}] \begin{bmatrix} 0\\0\\1 \end{bmatrix}$$





- **Answer:** All the components of the force and moment vectors are resisted by ٠ the structure of mechanism itself, except for the torque about the joint axis (revolute joint) or the force along the joint (prismatic joint).
- Therefore, to find the joint the torque or force required to maintain the static ٠ equilibrium, the dot product of the joint axis vector with the moment vector or force vector acting on the link is computed

Revolute Joint
$$\tau_{i} = {}^{i}n_{i}^{T}\hat{z}_{i} = [{}^{i}n_{ix} {}^{i}n_{iy} {}^{i}n_{iz}] \begin{bmatrix} 0\\0\\1 \end{bmatrix}$$
$$f_{i} = {}^{i}f_{i}^{T}\hat{z}_{i} = [{}^{i}f_{ix} {}^{i}f_{iy} {}^{i}f_{iz}] \begin{bmatrix} 0\\0\\1 \end{bmatrix}$$
Prismatic Joint

Revolute Joint



Solution

• Apply the static equilibrium equations starting from the end effector and going toward the base

Static force "propagation" from link to link

 ${}^{i}f_{i} = {}^{i}_{i+1}R {}^{i+1}f_{i+1}$ ${}^{i}n_{i} = {}^{i}_{i+1}R {}^{i+1}n_{i+1} + {}^{i}P_{i+1} \times {}^{i}f_{i}$

• To find the joint torque or force required to maintain the static equilibrium, the dot product of the joint axis vector with the moment vector or force vector acting on the link is computed

$$\tau_{i} = {}^{i}n_{i}^{T}\hat{z}_{i} = [{}^{i}n_{ix} {}^{i}n_{iy} {}^{i}n_{iz}] \begin{bmatrix} 0\\0\\1 \end{bmatrix} \qquad f_{i} = {}^{i}f_{i}^{T}\hat{z}_{i} = [{}^{i}f_{ix} {}^{i}f_{iy} {}^{i}f_{iz}] \begin{bmatrix} 0\\0\\1 \end{bmatrix}$$
Revolute Joint Prismatic Joint





SCARA – RRRP – DH Parameter (Modified form)



<i>i-1</i>	i	$lpha_{i-1}$	a_{i-1}	d_i	$ heta_i$
0	1	0	0	0	θ_1
1	2	0	l_1	0	θ_2
2	3	0	l_2	0	θ_3
3	4	180°	0	d_4	0



SCARA – RRRP

syms <mark>pi</mark>

```
L(1) = Link('revolute','d',0,'a',0,'alpha',0,'modified');
L(2) = Link('revolute','d',0,'a',l1,'alpha',0,'modified');
L(3) = Link('revolute','d',0,'a',l2,'alpha',0,'modified');
L(4) = Link('prismatic','alpha',pi,'theta', 0,'a',0,'modified')
```

```
SCARA = SerialLink(L, 'name', 'SCARA')
```

For plot [mm]

```
l1 = 0.3; l2=0.3; pi= 3.14;
t1 =0; t2=0; t3=0; d4=0.2;
th= [t1 t2 t3 d4];
L_P(1) = Link('revolute','d',0,'a',0,'alpha',0,'modified');
L_P(2) = Link('revolute','d',0,'a',11,'alpha',0,'modified');
L_P(3) = Link('revolute','d',0,'a',12,'alpha',0,'modified');
L_P(4) = Link('prismatic','alpha',pi,'theta', 0,'a',0,'modified')
SCARA_P = SerialLink(L_P,'name','SCARA')
figure()
SCARA_P.plot(th,'workspace',[-1 1 -1 1 -1 1])
```

SCARA:: 4 axis, RRRP, modDH, slowRNE

++					
ij	theta	d	a	alpha	offset
1	q1	0	0	0	0
2	q2	0	11	0	0
3	q3	0	12	0	0
4	0	q4	0	pi	0
++					



Simplify Function

 $T0_4 = SCARA.A([1 2 3 4],q)$ $T0_4 =$ $\begin{aligned} \cos(t_3) \,\sigma_2 - \sin(t_3) \,\sigma_3 & \sigma_1 & 0 \quad l_2 \,\sigma_2 + l_1 \cos(t_1) \\ \sigma_1 & \sin(t_3) \,\sigma_3 - \cos(t_3) \,\sigma_2 & 0 \quad l_2 \,\sigma_3 + l_1 \sin(t_1) \end{aligned}$ 0 0 -1 $-d_4$ 0 0 0 1 where $\sigma_1 = \cos(t_3) \,\sigma_3 + \sin(t_3) \,\sigma_2$ $\sigma_2 = \cos(t_1)\cos(t_2) - \sin(t_1)\sin(t_2)$ $\sigma_3 = \cos(t_1)\sin(t_2) + \cos(t_2)\sin(t_1)$

simplified_T0_4 = simplify(SCARA.A([1 2 3 4],q))

simplified_T0_4 =

1	$\cos(t_1 + t_2 + t_3)$	$\sin(t_1 + t_2 + t_3)$	0	$l_2 \cos(t_1 + t_2) + l_1 \cos(t_1)$
	$\sin(t_1 + t_2 + t_3)$	$-\cos(t_1+t_2+t_3)$	0	$l_2 \sin(t_1 + t_2) + l_1 \sin(t_1)$
l	0	0	-1	$-d_4$
(0	0	0	1 /



syms pi

3D-1-RRR Modi

```
clear all; close all;
syms t1 t2 t3 l2 l3
 L1(1) = Link('revolute', 'd',0, 'a',0, 'alpha',0, 'modified');
 L1(2) = Link('revolute','d',0,'a',0,'alpha',-pi/2,'modified');
 L1(3) = Link('revolute','d',0,'a',12,'alpha',0,'modified');
 L1(4) = Link('revolute','d',0,'a',13,'alpha',0,'modified');
 Robot_3D_1_RRR = SerialLink(L1, 'name', '3D-1-RRR Modi')
```

Robot 3D 1 RRR =



3D-1-RRR Modi

```
clear all; close all;
syms t1 t2 t3 l2 l3
svms pi
```

```
L1(1) = Link('revolute','d',0,'a',0,'alpha',0,'modified');
L1(2) = Link('revolute','d',0,'a',0,'alpha',-pi/2,'modified');
L1(3) = Link('revolute','d',0,'a',12,'alpha',0,'modified');
L1(4) = Link('revolute','d',0,'a',13,'alpha',0,'modified');
Robot_3D_1_RRR = SerialLink(L1, 'name', '3D-1-RRR Modi')
```

 $Robot_3D_1_RRR =$

 $\sigma_1 = l_3 \cos(t_2 + t_3) + l_2 \cos(t_2)$

3D-1	-RRR Modi:: 4	axis, RRRR,	modDH, slow	RNE	
j	theta	d	a	alpha	offset
1 2 3 4	q1 q2 q3 q4	0 0 0 0	0 0 12 13	0 -pi/2 0 0	0 0 0
3D-1-I T0_4 :	RRR Modi_Direct =	Kinematics			
	$(t_2 + t_3) \cos(t_1)$ $(t_2 + t_3) \sin(t_1)$ $-\sin(t_2 + t_3)$ 0	$-\sin(t_2 + t_3) = -\sin(t_2 + t_3) = -\cos(t_2 + t_3)$	$ \begin{aligned} \cos(t_1) & -\sin(t_1) \\ \sin(t_1) & \cos(t_1) \\ t_3) & 0 \\ & 0 \end{aligned} $) $\cos i \theta = -l_3 \sin(t_2 + \theta)$	$\begin{pmatrix} (t_1) \sigma_1 \\ (t_1) \sigma_1 \\ (t_3) - l_2 \sin(t_2) \\ 1 \end{pmatrix}$
where	2				



• The recursive equations for static force and torque from link *i*+1 to link *i*

$${}^{i}f_{i} = {}^{i}_{i+1}R^{i+1}f_{i+1}$$
$${}^{i}n_{i} = {}^{i}_{i+1}R^{i+1}n_{i+1} + {}^{i}P_{i+1} \times {}^{i}f_{i}$$

- It allows us to start with the force and torque applied at the end effector.
- By using this recursive equation, we can calculate the force and torque at each joint all the way back to the robot base frame





Force propagation

$${}^{i}f_{i} = {}_{i+1}{}^{i}R^{i+1}f_{i+1}$$

$${}^{4}f_{4} = \begin{bmatrix} f_{x} \\ f_{y} \\ f_{z} \end{bmatrix}$$

$${}^{3}f_{3} = {}_{4}{}^{3}R^{4}f_{4}$$

$${}^{2}f_{2} = {}_{3}{}^{2}R^{3}f_{3}$$

$${}^{1}f_{1} = {}_{2}{}^{1}R^{2}f_{2}$$

f4 = [fx;fy;fz] f3 = R3_4*f4; f3 = collect(f3,[fx fy fz]) f2 = R2_3*f3; f2 = collect(f2,[fx fy fz]) f1 = R1_2*f2; f1 = simplify(collect(f1,[fx fy fz]))



Force propagation

$${}^{i}f_{i} = {}_{i+1}{}^{i}R^{i+1}f_{i+1}$$

$${}^{4}f_{4} = \begin{bmatrix} f_{x} \\ f_{y} \\ f_{z} \end{bmatrix}$$

$${}^{3}f_{3} = {}_{4}^{3}R {}^{4}f_{4}$$

$${}^{2}f_{2} = {}_{3}^{2}R {}^{3}f_{3}$$

$${}^{1}f_{1} = {}_{2}^{1}R {}^{2}f_{2}$$

f4 = [fx;fy;fz] f3 = R3_4*f4; f3 = collect(f3,[fx fy fz]) f2 = R2_3*f3; f2 = collect(f2,[fx fy fz]) f1 = R1_2*f2; f1 = simplify(collect(f1,[fx fy fz])) $R0_{1} = t2r(T0_{1})$ $R1_{2} = t2r(T1_{2})$ $R2_{3} = t2r(T2_{3})$ $R3_{4} = t2r(T3_{4})$ $R0_{1} = \begin{pmatrix} \cos(t_{1}) & -\sin(t_{1}) & 0 \\ \sin(t_{1}) & \cos(t_{1}) & 0 \\ 0 & 0 & 1 \end{pmatrix}$ $R1_{2} = \begin{pmatrix} \cos(t_{2}) & -\sin(t_{2}) & 0 \\ \sin(t_{2}) & \cos(t_{2}) & 0 \end{pmatrix}$

 $\begin{pmatrix} \cos(t_2) & -\sin(t_2) \\ \sin(t_2) & \cos(t_2) \\ 0 & 0 \\ \text{R2_3} = \\ \begin{pmatrix} \cos(t_3) & -\sin(t_3) \\ \sin(t_3) & \cos(t_3) \\ 0 & 0 \\ \text{R3_4} = \\ \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ \end{pmatrix}$

0

 $T0_1 = SCARA.A([1],q)$ $T1_2 = SCARA.A([2],q)$ $T2_3 = SCARA.A([3],q)$ $T3_4 = SCARA.A([4],q)$ T0_1 = $\cos(t_1) - \sin(t_1) = 0$ $sin(t_1)$ $\cos(t_1) = 0 = 0$ 0 0 1 0 0 0 0 1 $T1_2 =$ $\cos(t_2) - \sin(t_2) = 0 \quad l_1$ $\sin(t_2) = \cos(t_2) = 0 = 0$ 1 0 0 0 0 0 0 1 T2_3 = $\cos(t_3) - \sin(t_3) = 0 \quad l_2$ $\sin(t_3) = \cos(t_3) = 0 = 0$ 0 0 1 0 0 0 0 1 $T3_4 =$ 0 -1 0 0



f4 =**Force propagation** ٠ ${}^{4}f_{4} = \begin{bmatrix} f_{x} \\ f_{y} \\ f_{z} \end{bmatrix}$ ${}^{3}f_{3} = {}^{3}_{4}R {}^{4}f_{4}$ ${}^{fy}_{fz}$ ${}^{f3} = {}^{(f)}_{4}$ ${}^{i}f_{i} = {}^{i}_{i+1}R^{i+1}f_{i+1}$ $\begin{pmatrix} fx \\ -fy \end{pmatrix}$ ${}^{2}f_{2} = {}^{2}_{3}R {}^{3}f_{3}$ f2 = ${}^{1}f_{1} = {}^{1}R {}^{2}f_{2}$ $\begin{pmatrix} \cos(t_3) \ fx + \sin(t_3) \ fy\\ \sin(t_3) \ fx + (-\cos(t_3)) \ fy\\ - fz \end{pmatrix}$ f4 = [fx;fy;fz] $f3 = R3_4 * f4;$ f3 = collect(f3, [fx fy fz])f1 = $f2 = R2_3 * f3;$ $\begin{pmatrix} fx \cos(t_2 + t_3) + fy \sin(t_2 + t_3) \\ fx \sin(t_2 + t_3) - fy \cos(t_2 + t_3) \\ - fz \end{pmatrix}$ f2 = collect(f2, [fx fy fz])f1 = R1 2*f2;f1 = simplify(collect(f1,[fx fy fz]))



• Torque propagation ${}^{i}n_{i} = {}^{i}_{i+1}R^{i+1}n_{i+1} + {}^{i}P_{i+1} \times {}^{i}f_{i}$

$${}^{4}n_{4} = \begin{bmatrix} n_{x} \\ n_{y} \\ n_{z} \end{bmatrix}$$

$${}^{3}n_{3} = {}^{3}_{4}R {}^{4}n_{4} + {}^{3}P_{4} \times {}^{3}f_{3}$$

$${}^{2}n_{2} = {}^{2}_{3}R {}^{3}n_{3} + {}^{2}P_{3} \times {}^{2}f_{2}$$

$${}^{1}n_{1} = {}^{1}_{2}R {}^{2}n_{2} + {}^{1}P_{2} \times {}^{1}f_{1}$$

n4 = [nx;ny;nz] n3 = R3_4*n4 + cross(P3_4,f3); n3 = collect(n3,[fx fy fz nx ny nz]) n2 = R2_3*n3 + cross(P2_3,f2); n2 = simplify(collect(n2,[fx fy fz nx ny nz])) n2 = collect(n2,[fx fy fz nx ny nz]) n1 = R1_2*n2 + cross(P1_2,f1);



- ${}^{i}n_{i} = {}^{i}_{i+1}R^{i+1}n_{i+1} + {}^{i}P_{i+1} \times {}^{i}f_{i}$ **Torque propagation** n4 = n_x nx $^{4}n_{4} = \mid n_{y} \mid$ ny nz n3 = ${}^{3}n_{3} = {}^{3}_{4}R {}^{4}n_{4} + {}^{3}P_{4} \times {}^{3}f_{3}$ n2 = ${}^{2}n_{2} = {}^{2}_{3}R {}^{3}n_{3} + {}^{2}P_{3} \times {}^{2}f_{2}$ $(d_4 \sin(t_3))$ fx + σ_1 fy + cos(t_3) nx + sin(t_3) ny $\sigma_1 \text{ fx} + (-d_4 \sin(t_3)) \text{ fy} + l_2 \text{ fz} + \sin(t_3) \text{ nx} + (-\cos(t_3)) \text{ ny}$ $(l_2 \sin(t_3)) \text{ fx} + (-l_2 \cos(t_3)) \text{ fy} - \text{nz}$ ${}^{1}n_{1} = {}^{1}_{2}R {}^{2}n_{2} + {}^{1}P_{2} \times {}^{1}f_{1}$ where $\sigma_1 = -d_4 \cos(t_3)$ n4 = [nx;ny;nz]n1 = $\begin{pmatrix} \sigma_2 \ fx + \sigma_1 \ fy + (-l_2 \sin(t_2)) \ fz + \cos(t_2 + t_3) \ nx + \sin(t_2 + t_3) \ ny \\ \sigma_1 \ fx + (-\sigma_2) \ fy + (l_1 + l_2 \cos(t_2)) \ fz + \sin(t_2 + t_3) \ nx + (-\cos(t_2 + t_3)) \ ny \\ (l_1 \sin(t_2 + t_3) + l_2 \sin(t_3)) \ fx + (-l_1 \cos(t_2 + t_3) - l_2 \cos(t_3)) \ fy - nz \end{pmatrix}$ $n3 = R3_4*n4 + cross(P3_4, f3);$ n3 = collect(n3, [fx fy fz nx ny nz]) $n2 = R2_3*n3 + cross(P2_3, f2);$ n2 = simplify(collect(n2, [fx fy fz nx ny nz]))n2 = collect(n2,[fx fy fz nx ny nz]) where n1 = R1 2*n2 + cross(P1 2, f1); $\sigma_1 = -d_4 \cos(t_2 + t_3)$
 - $\sigma_2 = d_4 \sin(t_2 + t_3)$



• To find the joint torque required to maintain the static equilibrium for a revolute joint,

$$\tau_i = {}^{i}n_i {}^{T}\hat{z}_i = \begin{bmatrix} {}^{i}n_{i,x} {}^{i}n_{i,y} {}^{i}n_{i,z} \end{bmatrix} \begin{bmatrix} 0\\0\\1 \end{bmatrix}$$

• To find the joint force required to maintain the static equilibrium for a prismatic joint,

$$f_{i} = {}^{i}f_{i}^{T}\hat{z}_{i} = \begin{bmatrix} {}^{i}f_{i,x} & {}^{i}f_{i,y} & {}^{i}f_{i,z} \end{bmatrix} \begin{bmatrix} 0\\0\\1 \end{bmatrix}$$

• We are only interested in ${}^{i}n_{i,z}$ and ${}^{i}f_{i,z}$

$$[\tau] = \begin{bmatrix} \tau_1 \\ \tau_2 \\ \tau_3 \\ \tau_4 \end{bmatrix} = \begin{bmatrix} {}^{1}n_{1,z} \\ {}^{2}n_{2,z} \\ {}^{3}n_{3,z} \\ {}^{4}f_{4,z} \end{bmatrix}$$

$$[\tau] = \begin{bmatrix} \tau_1 \\ \tau_2 \\ \tau_3 \\ \tau_4 \end{bmatrix} = \begin{bmatrix} {}^1n_{1,z} \\ {}^2n_{2,z} \\ {}^3n_{3,z} \\ {}^4f_{4,z} \end{bmatrix} = \begin{bmatrix} [{}^1n_{1,x} & {}^1n_{1,y} & {}^1n_{1z}] \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} \\ {}^2n_{2,z} \\ {}^3n_{3,z} \\ {}^3n_{3,z} \\ {}^4f_{4,x} & {}^4f_{4,y} & {}^4f_{4,z} \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$$

$$\begin{aligned} \mathsf{tq1} &= (l_1 \sin(t_2 + t_3) + l_2 \sin(t_3)) \ \mathsf{fx} + (-l_1 \cos(t_2 + t_3) - l_2 \cos(t_3)) \ \mathsf{fy} - \mathsf{nz} \\ \mathsf{tq2} &= (l_2 \sin(t_3)) \ \mathsf{fx} + (-l_2 \cos(t_3)) \ \mathsf{fy} - \mathsf{nz} \\ \mathsf{tq3} &= -\mathsf{nz} \end{aligned}$$

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tq4 = fz

UC

The Jacobian $\begin{pmatrix} 4 \\ J_4 \end{pmatrix}$ is defined as

$$[\tau] = \begin{bmatrix} {}^{4}J_{4} \end{bmatrix}^{T} \begin{bmatrix} f_{x} \\ f_{y} \\ f_{z} \\ n_{x} \\ n_{y} \\ n_{z} \end{bmatrix}$$

According to the definition
$$\begin{bmatrix} \tau_{1} \\ \tau_{2} \\ \tau_{3} \\ \tau_{4} \end{bmatrix} = \begin{bmatrix} {}^{1}n_{1,z} \\ {}^{2}n_{2,z} \\ {}^{3}n_{3,z} \\ {}^{4}f_{4,z} \end{bmatrix},$$
$$\begin{bmatrix} {}^{1}n_{1,z} \\ {}^{2}n_{2,z} \\ {}^{3}n_{3,z} \\ {}^{4}f_{4,z} \end{bmatrix} = \begin{bmatrix} {}^{4}J_{4} \end{bmatrix}^{T} \begin{bmatrix} f_{x} \\ f_{y} \\ f_{z} \\ n_{x} \\ n_{y} \\ n_{z} \end{bmatrix}$$

$$\begin{bmatrix} {}^{1}n_{1,z} \\ {}^{2}n_{2,z} \\ {}^{3}n_{3,z} \\ {}^{4}f_{4,z} \end{bmatrix} = \begin{bmatrix} {}^{4}J_{4} \end{bmatrix}^{T} \begin{bmatrix} f_{x} \\ f_{y} \\ f_{z} \\ n_{x} \\ n_{y} \\ n_{z} \end{bmatrix}$$

TQ1 = equationsToMatrix([tq1],[fx fy fz nx ny nz]) TQ2 = equationsToMatrix([tq2],[fx fy fz nx ny nz]) TQ3 = equationsToMatrix([tq3],[fx fy fz nx ny nz]) TQ4 = equationsToMatrix([tq4],[fx fy fz nx ny nz]) J4_FP_T = [TQ1;TQ2;TQ3;TQ4] J4_FP = transpose(J4_FP_T)

$$TQ1 = (l_1 \sin(t_2 + t_3) + l_2 \sin(t_3) - l_1 \cos(t_2 + t_3) - l_2 \cos(t_3) \ 0 \ 0 \ 0 \ -1)$$

$$TQ2 = (l_2 \sin(t_3) - l_2 \cos(t_3) \ 0 \ 0 \ 0 \ -1)$$

$$TQ3 = (0 \ 0 \ 0 \ 0 \ 0 \ -1)$$

$$TQ4 = (0 \ 0 \ 1 \ 0 \ 0 \ 0)$$



$$\begin{bmatrix} 4 J_4 \end{bmatrix} = \begin{bmatrix} 4 J_4 \end{bmatrix}^T$$

J4_FP_T = [TQ1;TQ2;TQ3;TQ4]
J4_FP = transpose(J4_FP_T)







Jacobian Methods of Derivation & the Corresponding Reference Frame – Summary

Method	Jacobian Matrix Reference Frame	Transformation	on to Base Frame (Frame 0)
Explicit (Diff. the Forward Kinematic Eq.)	${}^{0}{J}_{_{N}}$	None	
Iterative Velocity Eq.	$^{\scriptscriptstyle N}J_{\scriptscriptstyle N}$	Transform Method 1: ${}^{0}v$ ${}^{0}\alpha$ Transform Method 2: ${}^{0}J_{j}$	$\mathcal{V}_{N} = {}_{N}^{0} R^{N} \mathcal{V}_{N}$ $\mathcal{D}_{N} = {}_{N}^{0} R^{N} \mathcal{D}_{N}$ $\mathcal{U}_{N}(\theta) = {} {}_{N}^{0} R^{N} \theta \\ 0 \qquad {}_{N}^{0} R \end{bmatrix} {}_{N} \mathcal{J}_{N}(\theta)$
Iterative Force Eq.	$^{\scriptscriptstyle N} J_{\scriptscriptstyle N}^{\scriptscriptstyle T}$	Transpose N_{J} Transform ${}^{0}J_{J}$	$J_{N} = \begin{bmatrix} {}^{N}J_{N}^{T} \end{bmatrix}^{T}$ $J_{N}(\theta) = \begin{bmatrix} {}^{0}R & 0 \\ 0 & {}^{0}R \end{bmatrix} {}^{N}J_{N}(\theta)$



Jacobian: Frame of Representation

% convert 4J4_FP to 0J4_FP
J0_FP=[R0_4 zeros(3,3); zeros(3,3) R0_4]*J4_FP;
J0_FP=simplify(J0_FP)



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SCARA example _ Jacobian Matrix

Differentiation the Forward Kinematics Eqs. (Method 1)	Velocity Propagation – Base to EE (Method 2)	Force/Torque Propagation – EE to Base (Method 3)
$\boldsymbol{J} = \begin{bmatrix} \boldsymbol{J}_P \\ \boldsymbol{J}_O \end{bmatrix}$		$\begin{array}{c} \exists 4_FP_T = \\ \begin{pmatrix} l_1 \sin(t_2 + t_3) + l_2 \sin(t_3) & -l_1 \cos(t_2 + t_3) - l_2 \cos(t_3) & 0 & 0 & 0 & -1 \\ l_2 \sin(t_3) & -l_2 \cos(t_3) & 0 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 & 0 & 0 \end{array} \right)$
	$\begin{array}{rllllllllllllllllllllllllllllllllllll$	$\begin{array}{c} D4_FP = \\ \begin{pmatrix} l_1 \sin(t_2 + t_3) + l_2 \sin(t_3) & l_2 \sin(t_3) & 0 & 0 \\ -l_1 \cos(t_2 + t_3) - l_2 \cos(t_3) & -l_2 \cos(t_3) & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ -1 & -1 & -1 & -1 & 0 \end{array}$
$ \begin{array}{c} \hline \textbf{D0}_\textbf{DD} = \\ \begin{pmatrix} -l_2 \sin(t_1 + t_2) - l_1 \sin(t_1) & -l_2 \sin(t_1 + t_2) & 0 & 0 \\ l_2 \cos(t_1 + t_2) + l_1 \cos(t_1) & l_2 \cos(t_1 + t_2) & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 1 & 1 & 1 & 0 \\ \end{array} $	$\begin{array}{l} J0_VP = \\ \left(\begin{array}{ccc} -l_2\sin(t_1+t_2) - l_1\sin(t_1) & -l_2\sin(t_1+t_2) & 0 & 0 \\ l_2\cos(t_1+t_2) + l_1\cos(t_1) & l_2\cos(t_1+t_2) & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 1 & 1 & 1 & 0 \end{array} \right) \end{array}$	$\begin{array}{c} J0_FP = \\ \begin{pmatrix} -l_2 \sin(t_1 + t_2) - l_1 \sin(t_1) & -l_2 \sin(t_1 + t_2) & 0 & 0 \\ l_2 \cos(t_1 + t_2) + l_1 \cos(t_1) & l_2 \cos(t_1 + t_2) & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 1 & 1 & 1 & 0 \end{array} \right)$

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Kinematics Relations - Forward & Inverse

• The robot kinematic equations relate the two description of the robot tip location







• The velocity relationship: The relationship between the joint angle rates (\dot{X}) and the linear and angular velocities of the end effector $(\dot{\theta})$.

 $\dot{\mathbf{v}} = \mathbf{i}\dot{\mathbf{o}}$

- The velocity relationship
 - : The relationship between the joint angle rates ($\underline{\dot{\theta}}_N$) and the translation and rotation velocities of the end effector ($\underline{\dot{x}}$).

$$\underline{\dot{x}} = J(\underline{\theta})\underline{\dot{\theta}}$$

 The relationship between the robot joint torques (<u>τ</u>) and the forces and moments (<u>F</u>) at the robot end effector (Static Conditions).

$$\underline{\tau} = J(\underline{\theta})^T \underline{F}$$





Jacobian Matrix - Derivation Methods







$$\dot{p}_e = J_P(q)\dot{q} \tag{3.2}$$

$$\omega_e = J_O(q)\dot{q}. \tag{3.3}$$

In (3.2) J_P is the $(3 \times n)$ matrix relating the contribution of the joint velocities \dot{q} to the end-effector *linear* velocity \dot{p}_e , while in (3.3) J_O is the $(3 \times n)$ matrix relating the contribution of the joint velocities \dot{q} to the end-effector *angular* velocity ω_e . In compact form, (3.2), (3.3) can be written as

$$v_e = \begin{bmatrix} \dot{p}_e \\ \omega_e \end{bmatrix} = J(q)\dot{q} \tag{3.4}$$

which represents the manipulator differential kinematics equation. The $(6 \times n)$ matrix J is the manipulator geometric Jacobian

$$J = \begin{bmatrix} J_P \\ J_O \end{bmatrix},\tag{3.5}$$

which in general is a function of the joint variables.

In summary, the Jacobian in (3.5) can be partitioned into the (3×1) column vectors j_{Pi} and j_{Oi} as

$$J = \begin{bmatrix} \mathcal{I}_{P1} & \mathcal{I}_{Pn} \\ & \dots & \\ \mathcal{I}_{O1} & \mathcal{I}_{On} \end{bmatrix}, \qquad (3.29)$$

where

$$\begin{bmatrix} \jmath_{Pi} \\ \jmath_{Oi} \end{bmatrix} = \begin{cases} \begin{bmatrix} z_{i-1} \\ 0 \end{bmatrix} & \text{for a prismatic joint} \\ \begin{bmatrix} z_{i-1} \times (p_e - p_{i-1}) \\ z_{i-1} \end{bmatrix} & \text{for a revolute joint.} \end{cases}$$
(3.30)

The expressions in (3.30) allow Jacobian computation in a simple, systematic way on the basis of direct kinematics relations. In fact, the vectors z_{i-1} , p_e and p_{i-1} are all functions of the joint variables. In particular:

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Jacobian Matrix - Derivation Methods







Velocity Propagation

- Given: A manipulator A chain of rigid bodies each one capable of moving relative to its neighbor
- Problem: Calculate the linear and angular velocities of the link of a robot
- Solution (Concept): Due to the robot structure (serial mechanism) we can compute the velocities of each link in order starting from the base.



The velocity of link *i*+1

- = The velocity of link *i*
 - + whatever new velocity components were added by joint *i*+1





Velocity Propagation – Intuitive Explanation

- Three Actions
 - The origin of frame B moves with respect to the origin of frame A
 - Point Q moves with respect to frame B
 - Frame B rotates with respect to frame A about an axis defined by ${}^{A}\Omega_{B}$

Vector Form

$${}^{A}V_{Q} = {}^{A}V_{BORG} + {}^{A}_{B}R^{B}V_{Q} + {}^{A}\Omega_{B} \times {}^{A}_{B}R^{B}P_{Q}$$

Matrix Form

$${}^{A}V_{Q} = {}^{A}V_{BORG} + {}^{A}_{B}R^{B}V_{Q} + {}^{A}_{B}\dot{R}_{\Omega}\left({}^{A}_{B}R^{B}P_{Q}\right)$$





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• The recursive equation for the Angular Velocity is

$$^{i+1}\omega_{i+1} = {}^{i+1}_{i}R {}^{i}\omega_{i} + \rho \begin{bmatrix} 0\\0\\\dot{\theta}_{i+1} \end{bmatrix}, \rho = 0 \text{ in the prismatic joint}$$

 $^{i+1}\rho = \mathbf{1} \text{ in the revolute joint}$

• The recursive equation for Linear Velocity is

$${}^{i+1}v_{i+1} = {}^{i+1}{}_{i}R\left({}^{i}\omega_{i} \times {}^{i}P_{i+1} + {}^{i}v_{i}\right) + \rho \begin{bmatrix} 0\\0\\d \end{bmatrix}, \rho = 1 \text{ in the prismatic joint} \\ \rho = \mathbf{0} \text{ in the revolute joint} \\ \rho = \mathbf{0} \text{ in the revolute joint} \\ {}^{4}v_{4,x} \\ {}^{4}v_{4,y} \\ {}^{4}v_{4,z} \\ {}^{4}w_{4,1} \\ {}^{4}w_{4,2} \\ {}^{4}w_{4,3} \end{bmatrix}$$



Jacobian Matrix - Derivation Methods





Solution

Apply the static equilibrium equations starting from the end effector and going toward the base

Static force "propagation" from link to link

$${}^{i}f_{i} = {}^{i}_{i+1}R {}^{i+1}f_{i+1}$$
$${}^{i}n_{i} = {}^{i}_{i+1}R {}^{i+1}n_{i+1} + {}^{i}P_{i+1} \times {}^{i}f_{i}$$

• To find **the joint torque or force required to maintain the static equilibrium**, the dot product of the joint axis vector with the moment vector or force vector acting on the link is computed

$$\tau_{i} = {}^{i}n_{i}^{T}\hat{z}_{i} = [{}^{i}n_{ix} \quad {}^{i}n_{iy} \quad {}^{i}n_{iz}] \begin{bmatrix} 0\\0\\1 \end{bmatrix} \qquad f_{i} = {}^{i}f_{i}^{T}\hat{z}_{i} = [{}^{i}f_{ix} \quad {}^{i}f_{iy} \quad {}^{i}f_{iz}] \begin{bmatrix} 0\\0\\1 \end{bmatrix}$$
Revolute Joint Prismatic Joint



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The transpose of Jacobian $\begin{pmatrix} 4 J_4^T \end{pmatrix}$ is defined as

$$[\tau] = \begin{bmatrix} {}^{4}J_{4} \end{bmatrix}^{T} \begin{bmatrix} J_{x} \\ f_{y} \\ f_{z} \\ n_{x} \\ n_{y} \\ n_{z} \end{bmatrix}$$

According to the definition
$$\begin{bmatrix} \tau_{1} \\ \tau_{2} \\ \tau_{3} \\ \tau_{4} \end{bmatrix} = \begin{bmatrix} {}^{1}n_{1,z} \\ {}^{2}n_{2,z} \\ {}^{3}n_{3,z} \\ {}^{4}f_{4,z} \end{bmatrix},$$
$$\begin{bmatrix} {}^{1}n_{1,z} \\ {}^{2}n_{2,z} \\ {}^{3}n_{3,z} \\ {}^{4}f_{4,z} \end{bmatrix} = \begin{bmatrix} {}^{4}J_{4} \end{bmatrix}^{T} \begin{bmatrix} f_{x} \\ f_{y} \\ f_{z} \\ n_{x} \\ n_{y} \\ n_{z} \end{bmatrix}$$

SCARA example _ Jacobian Matrix

Differentiation the Forward Kinematics Eqs. (Method 1)	Velocity Propagation – Base to EE (Method 2)	Force/Torque Propagation – EE to Base (Method 3)
$\boldsymbol{J} = \begin{bmatrix} \boldsymbol{J}_P \\ \boldsymbol{J}_O \end{bmatrix}$		$\begin{array}{c} \exists 4_FP_T = \\ \begin{pmatrix} l_1 \sin(t_2 + t_3) + l_2 \sin(t_3) & -l_1 \cos(t_2 + t_3) - l_2 \cos(t_3) & 0 & 0 & 0 & -1 \\ l_2 \sin(t_3) & -l_2 \cos(t_3) & 0 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 & 0 & 0 \end{array} \right)$
	$\begin{array}{rllllllllllllllllllllllllllllllllllll$	$\begin{array}{c} D4_FP = \\ \begin{pmatrix} l_1 \sin(t_2 + t_3) + l_2 \sin(t_3) & l_2 \sin(t_3) & 0 & 0 \\ -l_1 \cos(t_2 + t_3) - l_2 \cos(t_3) & -l_2 \cos(t_3) & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ -1 & -1 & -1 & -1 & 0 \end{array}$
$ \begin{array}{c} \hline \textbf{D0}_\textbf{DD} = \\ \begin{pmatrix} -l_2 \sin(t_1 + t_2) - l_1 \sin(t_1) & -l_2 \sin(t_1 + t_2) & 0 & 0 \\ l_2 \cos(t_1 + t_2) + l_1 \cos(t_1) & l_2 \cos(t_1 + t_2) & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 1 & 1 & 1 & 0 \\ \end{array} $	$\begin{array}{l} J0_VP = \\ \left(\begin{array}{ccc} -l_2\sin(t_1+t_2) - l_1\sin(t_1) & -l_2\sin(t_1+t_2) & 0 & 0 \\ l_2\cos(t_1+t_2) + l_1\cos(t_1) & l_2\cos(t_1+t_2) & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 1 & 1 & 1 & 0 \end{array} \right) \end{array}$	$\begin{array}{c} J0_FP = \\ \begin{pmatrix} -l_2 \sin(t_1 + t_2) - l_1 \sin(t_1) & -l_2 \sin(t_1 + t_2) & 0 & 0 \\ l_2 \cos(t_1 + t_2) + l_1 \cos(t_1) & l_2 \cos(t_1 + t_2) & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 1 & 1 & 1 & 0 \end{array} \right)$

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Jacobian Methods of Derivation & the Corresponding Reference Frame – Summary

Method	Jacobian Matrix Reference Frame	Transformation to Base Frame (Frame 0)
Explicit (Diff. the Forward Kinematic Eq.)	$^{0}{oldsymbol{J}}_{N}$	None
Iterative Velocity Eq.	$^{\scriptscriptstyle N}J_{\scriptscriptstyle N}$	Transform Method 1: ${}^{0}\nu_{N} = {}^{0}_{N}R^{N}\nu_{N}$ ${}^{0}\omega_{N} = {}^{0}_{N}R^{N}\omega_{N}$ Transform Method 2: ${}^{0}J_{N}(\theta) = \left[{}^{0}_{N}R 0 \\ 0 {}^{0}_{N}R \right] {}^{N}J_{N}(\theta)$
Iterative Force Eq.	$^{\scriptscriptstyle N}J_{\scriptscriptstyle N}^{\scriptscriptstyle T}$	Transpose ${}^{N}J_{N} = [{}^{N}J_{N}^{T}]^{T}$ Transform ${}^{0}J_{N}(\theta) = \begin{bmatrix} {}^{0}R & 0 \\ 0 & {}^{0}_{N}R \end{bmatrix} {}^{N}J_{N}(\theta)$

Summary

- ✓ Jacobian with SCARA example
 - Velocity propagation
 - Direct differentiation
 - Force/Torque propagation
- ✓ Frame of Representation

Next Discussion Section

Manipulator Dynamics - Concept

- Forward Dynamics
- Inverse Dynamics
- □ Manipulator Dynamics Solution
 - Newton-Euler Equations
 - Lagrangian Dynamics

Or Lab section with DENSO robot

