

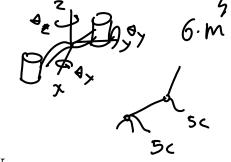
Manipulator Dynamics 4





1. Define a set of *generalized coordinates* for i=1,2,3...N. These coordinates can be chosen arbitrarily as long as they provide a set of poF independent variables that map the system in a 1-to-1 manner. The usual variable set for serial manipulators is:

 $q_i = \begin{cases} \theta_i \text{ if revolute joint} \\ d_i \text{ if prismatic joint} \end{cases}$



- **2.** Define a set of **generalized velocities** \dot{q}_i for *i=1,2,3...N*
- **3.** Define a set of *generalized forces (and moments)* Q_i for *i=1,2,3...N* The generalized forces must satisfy

$$Q_i \delta q_i = \delta W$$

where δq_i is a small change in the generalized coordinate and δW is the work done corresponding to that small change.

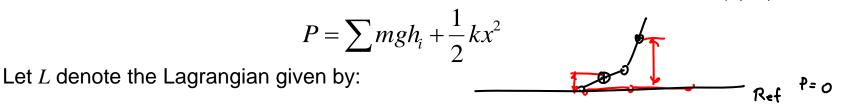


4. Write the equations describing the *kinetic and potential energies* as functions of the generalized coordinates as well as the resulting Lagrangian.

Let *K* denote the expression describing the kinetic energy. where $K = f(q_i, \dot{q}_i, t)$

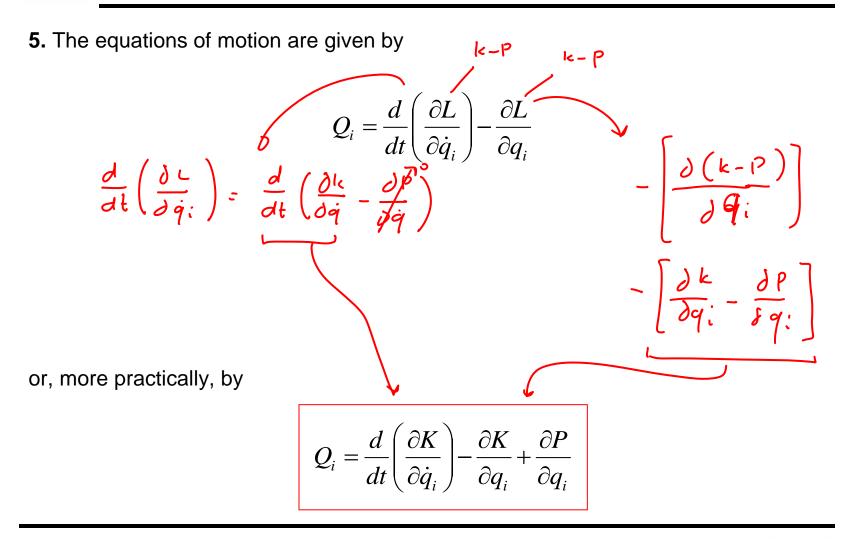
$$k_{i} = \frac{1}{2} v_{ci}^{T} m_{i} v_{ci} + \frac{1}{2} \omega_{i}^{ci} I_{i}^{i} \omega_{i}$$

Let *P* denote the expression describing the potential energy. where $P = f(q_i, t)$

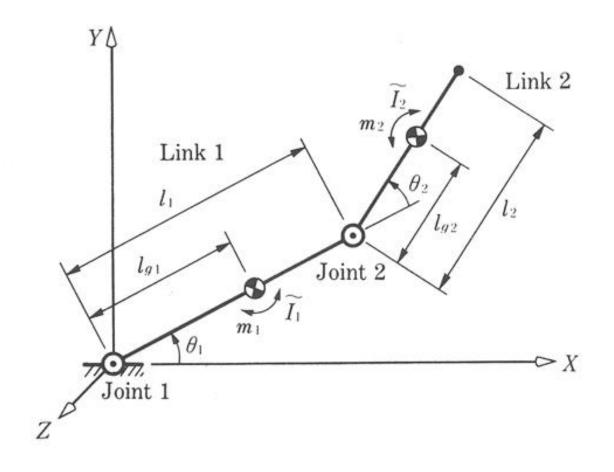


$$L=K-P$$

Langrangian Formulation of Manipulator Dynamics 3/









Step 1: Let
$$q_1 = \theta_1$$
 and $q_2 = \theta_2$

Step 2: Let $\dot{q}_1 = \dot{\theta}_1$ and $\dot{q}_2 = \dot{\theta}_2$

Step 3: Let external forces/torques $Q_i = \tau_i$

Step 4:

• Kinetic Energy:
$$k_i = \frac{1}{2} m_i v_{ci}^T v_{ci} + \frac{1}{2} \omega_i^{ci} I_i^i \omega_i$$

• For
$$i=1$$

 \longrightarrow $k_1 = \frac{1}{2}m_1L_{g1}^2\dot{\theta}_1^2 + \frac{1}{2}I_1\dot{\theta}_1^2$

$$V_{c_1} = W_{XF}$$

 $\dot{\forall} \times l_{g_1}$

UCI



 To find the velocity of the center of mass of link 2, first consider its position given by

$${}^{0}P_{g2} = \begin{bmatrix} L_{1}c1 + L_{g2}c12 \\ L_{1}s1 + L_{g2}s12 \end{bmatrix} = {}^{0}V_{c_{1}} = \begin{bmatrix} -L_{1}s(\dot{\theta}_{1} - L_{g2}s12) (\dot{\theta}_{1}+\dot{\theta}_{2}) (\dot{\theta}_{1}+\dot$$



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- Potential Energy: $p = \sum mgh_i$
- For i=1 $p_1 = m_1 g L_{g1} s 1$
- For i=2 $p_2 = m_2 g (L_1 s 1 + L_{g2} s 12)$

• Lagrangian: $L = k_1 + k_2 - p_1 - p_2$



• Step 5: Solving

$$Q_i = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i}$$

$$\longrightarrow Q_i = \frac{d}{dt} \left(\frac{\partial K}{\partial \dot{q}_i} \right) - \frac{\partial K}{\partial q_i} + \frac{\partial P}{\partial q_i}$$

$$\tau_{1} = \frac{d}{dt} \left(\frac{\partial K}{\partial \dot{\theta}_{1}} \right) - \frac{\partial K}{\partial \theta_{1}} + \frac{\partial P}{\partial \theta_{1}}$$
$$\tau_{2} = \frac{d}{dt} \left(\frac{\partial K}{\partial \dot{\theta}_{2}} \right) - \frac{\partial K}{\partial \theta_{2}} + \frac{\partial P}{\partial \theta_{2}}$$



$$\begin{aligned} \tau_{1} &= [m_{1}L_{g1} + I_{1} + m_{2}(L_{1}^{2} + L_{g2}^{2} + 2L_{1}L_{g2}c_{2}) + I_{2}]\ddot{\theta}_{1} \\ &+ [m_{2}(L_{g2}^{2} + L_{1}L_{g2}c_{2} + I_{2}]\ddot{\theta}_{2} \\ &- m_{2}L_{1}L_{g2}s_{2}(2\dot{\theta}_{1}\dot{\theta}_{2} + \dot{\theta}_{2}^{2}) \\ &+ m_{1}gL_{g1}c_{1} + m_{2}g(L_{1}c_{1} + L_{g2}c_{12}) \end{aligned}$$

$$\tau_{2} = [m_{2}(L_{g2}^{2} + L_{1}L_{g2}c_{2}) + I_{2}]\ddot{\theta}_{1}$$
$$+ [m_{2}L_{g2} + I_{2}]\ddot{\theta}_{2}$$
$$+ m_{2}L_{1}L_{g2}s_{2}\dot{\theta}_{1}^{2}$$
$$+ m_{2}gL_{g2}c_{12}$$



Gravity Effects - Langrangian Formulation

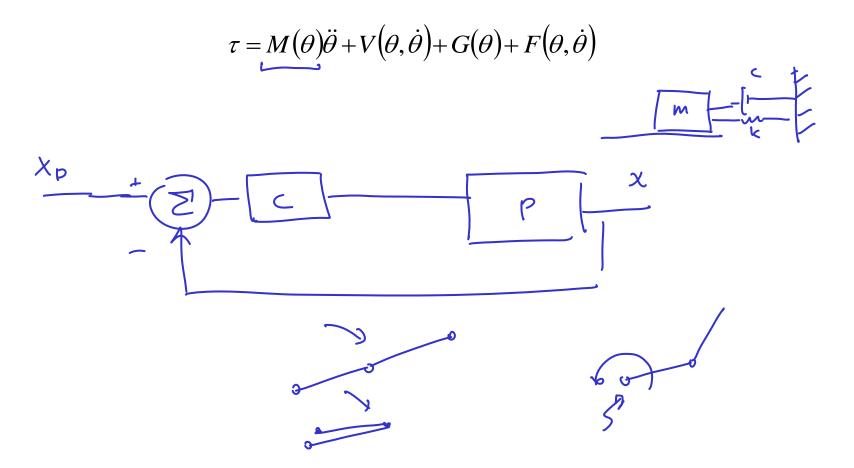
$$\tau_{i} = \frac{d}{dt} \left(\frac{\partial K(\theta, \dot{\theta})}{\partial \dot{\theta}_{i}} \right) - \frac{\partial K(\theta, \dot{\theta})}{\partial \theta_{i}} + \frac{\partial P(\theta)}{\partial \theta_{i}}$$

$$\tau = M(\theta)\ddot{\theta} + V(\theta,\dot{\theta}) + G(\theta)$$



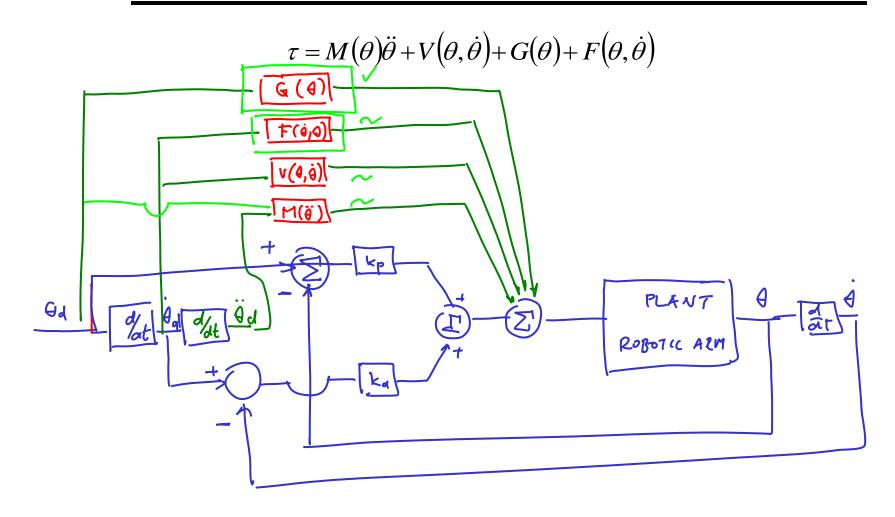


Manipulators - Control Problem





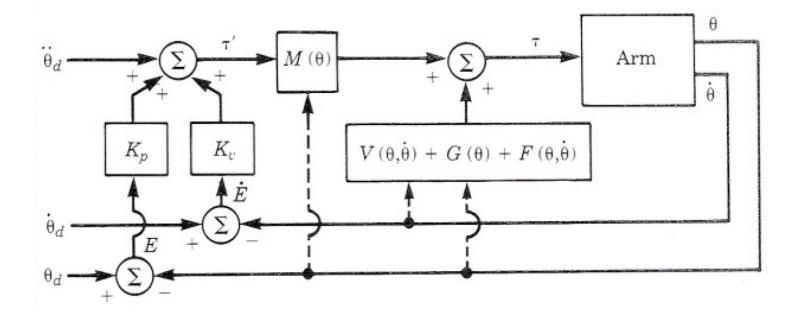
Manipulators – Non Linear Control Problem





Manipulators – Non Linear Control Problem

 $\tau = M(\theta)\ddot{\theta} + V(\theta, \dot{\theta}) + G(\theta) + F(\theta, \dot{\theta})$





Equation of Motion – Non Rigid Body Effects

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$$\tau = M(\theta)\ddot{\theta} + V(\theta, \dot{\theta}) + G(\theta) + F(\theta, \dot{\theta})$$

• Viscous Friction

$$\tau_{friction} = v\dot{\theta}$$

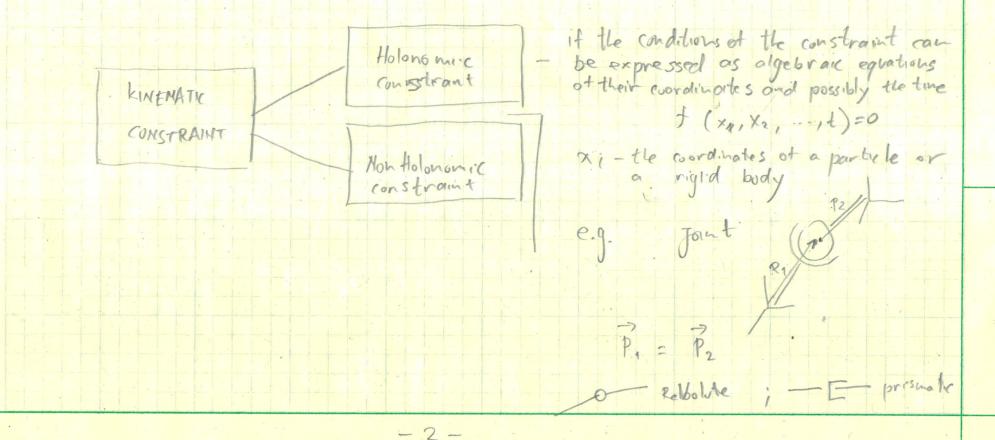
- Coulomb Friction $\tau_{friction} = c \operatorname{sgn}(\dot{\theta})$
- Model of Friction $\tau_{friction} = v\dot{\theta} + c \operatorname{sgn}(\dot{\theta}) = f(\theta, \dot{\theta})$

LAGRANGIAN FORMULATION OF MANIPULATOR DYNAMICS Lagrangian function - The difference between kinetic and potential engergy of a mechanical system. L= K - 1 L- Lagrangian K- kinetic engergy of a mechanical system II- Potential engergy of a mechanical system kinetic energy K=f(P,V) - function of the position & velocity of the potential energy U=f(P) - time tion of the postson of the link lagrange's equation of motion is $\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{q}}\right) - \frac{\partial L}{\partial \dot{q}_i} = Qi$



9 - vector of generlaed coordinates 9=[9,92, ...9n] Q - Vector of generlized forces Q = [Q1, Q2 - Qn]

Kinematic Constraint - Imposes some condition on the relative motion between a pair of bodies (e.g. - joint)





- The configuration of a mechanical system is known completely it the position and orientation of all the bodies in the system with respect to a reference frame a known Rigid body 6 DOF ×, 1, 2, 8x, 87, 62 Mechanical system with (M) moving bodies requires (6 m) coordinates to specify its configuration completely in 3D space - Robotic Arm - link are subject to mechanical constrains imposed by the joint the (6m) coordinates are no longer independent Most of the constrains encountered in a robotic system are holohomix



C - No of holonomic constrains h - DOF h = 6M - CThe (n) independent variables are a set of independent generalized coordinates # independent genoralized = [# DOF] coordinates * Lagrangian Equations of the second Type - All the forces of constraint in the joint do not appear in the equations, # of eq = # 00F Applicable to mechanical system with holonamic constraints Applicable to serial manipulators

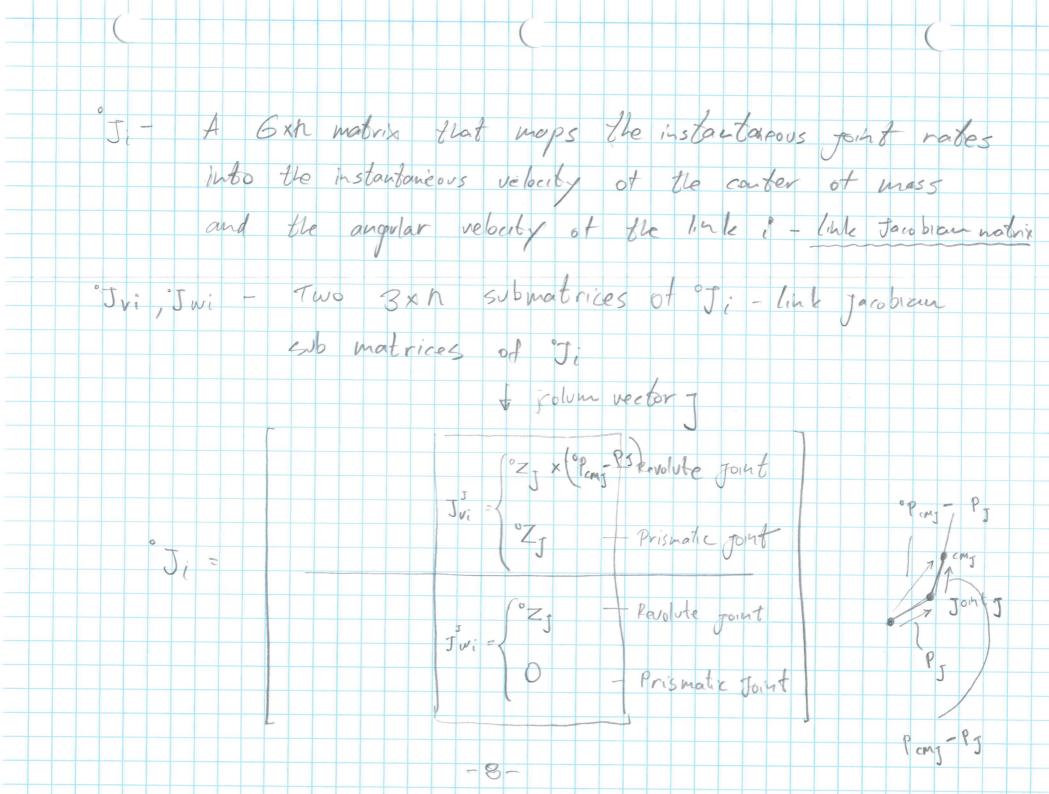


* Lagrangian Equations of the First Type - # nonindependent coordinates > # DOF - coordinates ave no longer independent [Eq of motion L'set of constraint Eq. (e.g. Lagrangian Multipliers) nonindependent coordinates - Redundant Lagrangian Coordinates - sutable for paralles manipulators - Applicable to mechanical systems with both holonamic and nonholonamic constraints FOR SERIAL MANIPULATOR Joint angle (revolute joint) 9=[91,92---9n] .9 translational distance (prismatic joint)

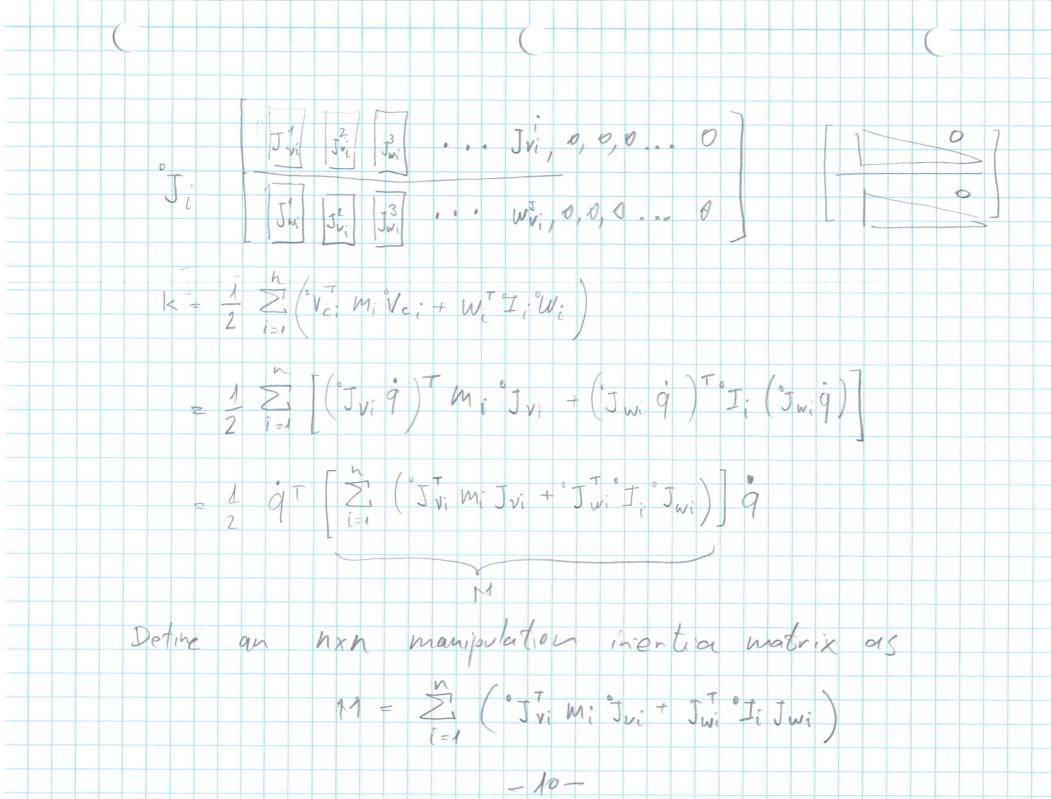
DIMENSIONS 9: (RAD) [m] Qiq: [J] $Q_i = [N_m]$ EINETIC ENERGY link i 0 Ki = 1 Vci Mi Vci + 1 Wi Jie Wi "I:c = The inertia matrix of link i about its CM and expressed in the base frame 'Iie -1 - While frame



Ficm = Ri Fice (Ri)T I ic - Time Invariant NOTE : "Tie - Depends on the rabot and posture because it is expressed in the base trave and the orientation of link i with respect to the base is a function of found variables Mathods for expressing the relacity at the center of mass "Ve: - Recursive method Jarobian (instantaneous screw notion) Vei = Ji Ai $v_{ci} = V_{ci}$ $J = J_{vi}$



Pci - Position veter defined from the orgine of the g link frame to the conter of mass of links i and expressed in ble base trave Jui, Jui - The pandral rate of change of the relocity of the center of mass and the angular velocity of link i with respect to the Joint Jth jout motion Since the motion of link i depends only on the point 1 through i de two vectors in the matrix are set to zero for J>i wind wi Depends on the previous joints -9-



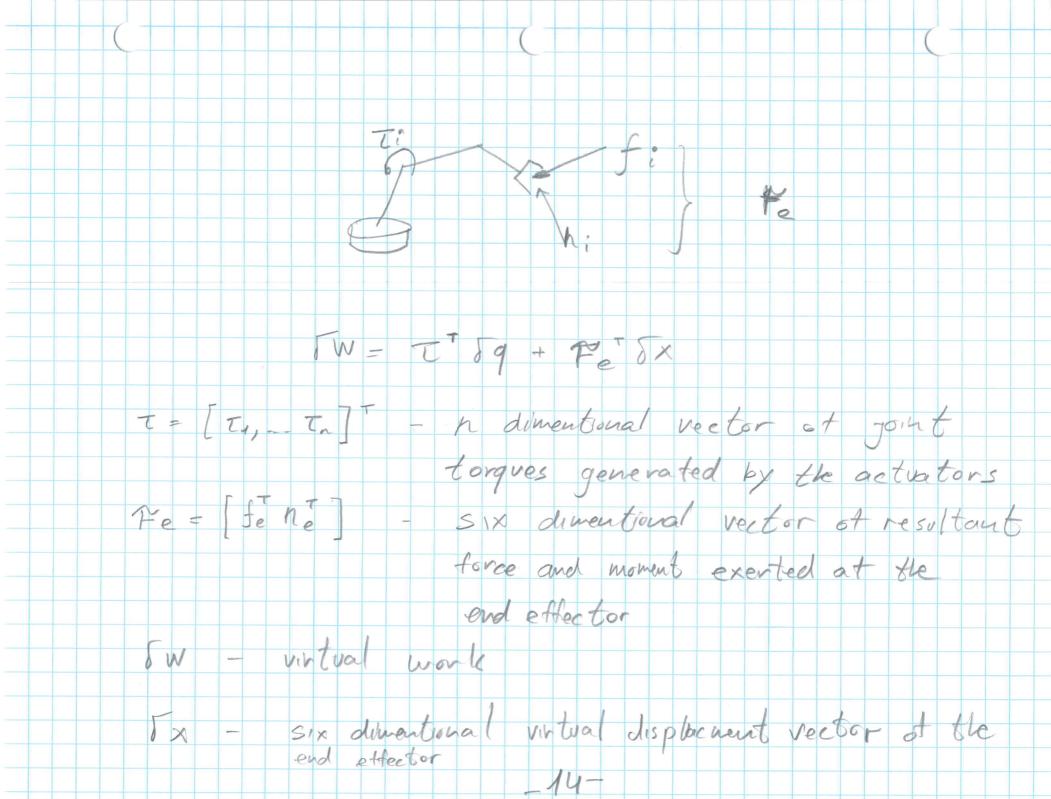
The total kinetic energy of a robot arm can be expressed in terms of the manipulator inertia matrix and the vector of joint vales $|K = \frac{1}{2} q^T M q$ Mis configuration dependent because Judju ane contiguration dependent as well M - symatuc - positive definite The guadratic form of the equation indicates that the kinetic energy is always positive unless the system is at rest.

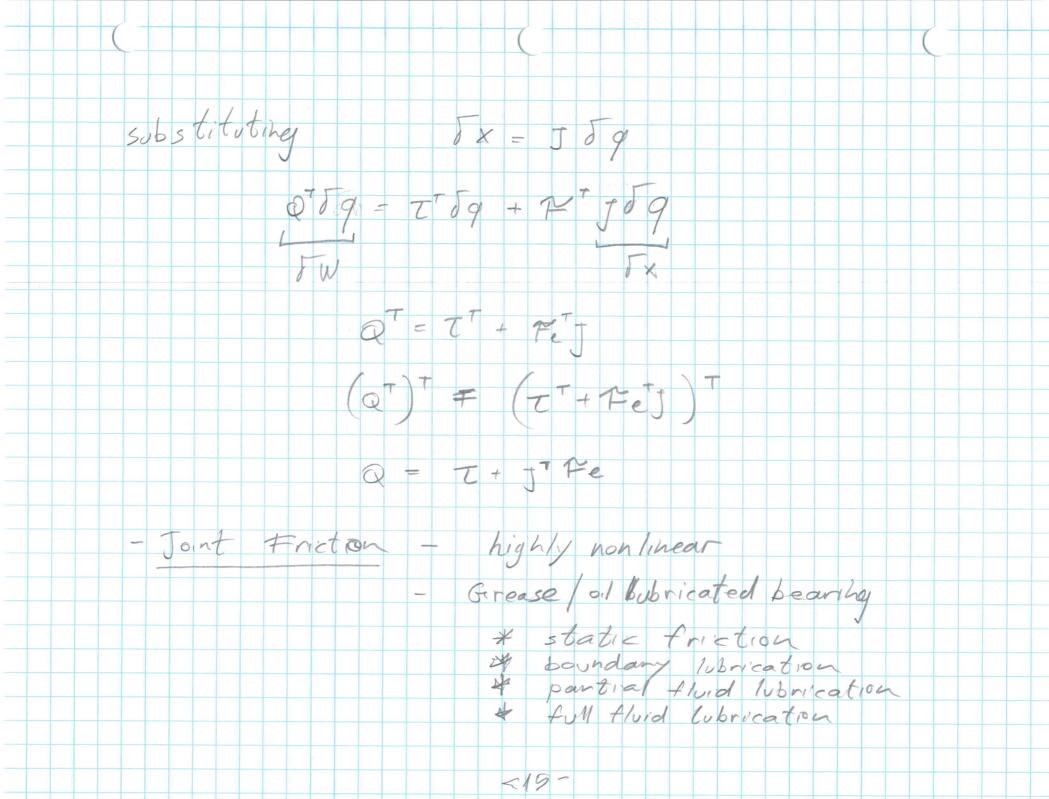
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POTENTIAL EXERGY

POTENTIAL EXERGY stored in a link is defined as the amout at work required to raise the center at mass of link i from the horizontal neterance plane to it present position under the influence of gravity with reference to the inertial frame (frame O) the work required to displace link (i) to position pci is given by Mig Pci. The total potential energy stored in a robot arm is $T = -\sum_{i=1}^{2} m_i g P_{ci}$ × -12-

FORCES GENERALIZED Gravitation forces FORCES Intertial forces All the rest - Growtation forces All the forces acting on a robot and that consistent with the mechanical constraints The vector of generalized forces Q = [Q1, Q2, ... Qn] is defined by the principle of vintual work as FW = aT FqActuators -> Force Torque at the forht External Forces/Moment -> END EPPECTOR -13-

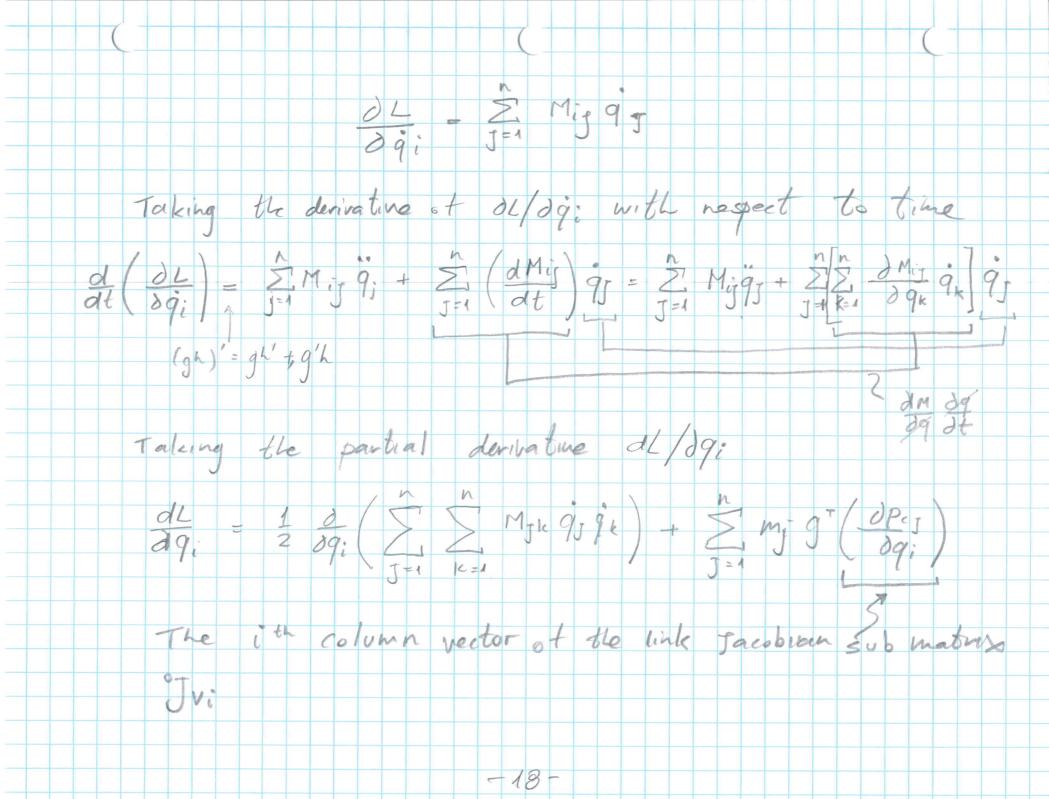




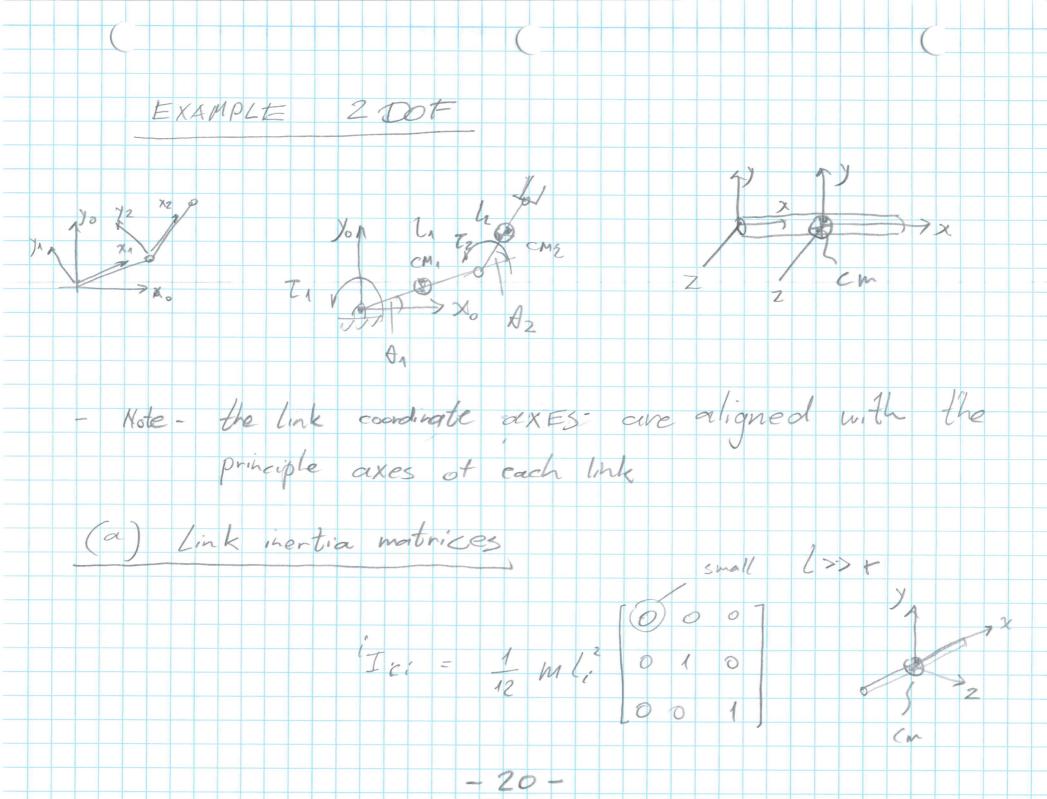
- Full fluid Lubrication Frication force & relative vebcrty $f_{i} = -b_{i} q_{i}$ The virtual work contributed by this type of frication $\delta w = -f_F \delta q$ fr = | bigi, bigi - on gn] - The frictional torques or forces in the joints. The minus sign indicates that the direction of the fricational targue or force is always opposite to the joint veherby R= t+ j Fe - fr

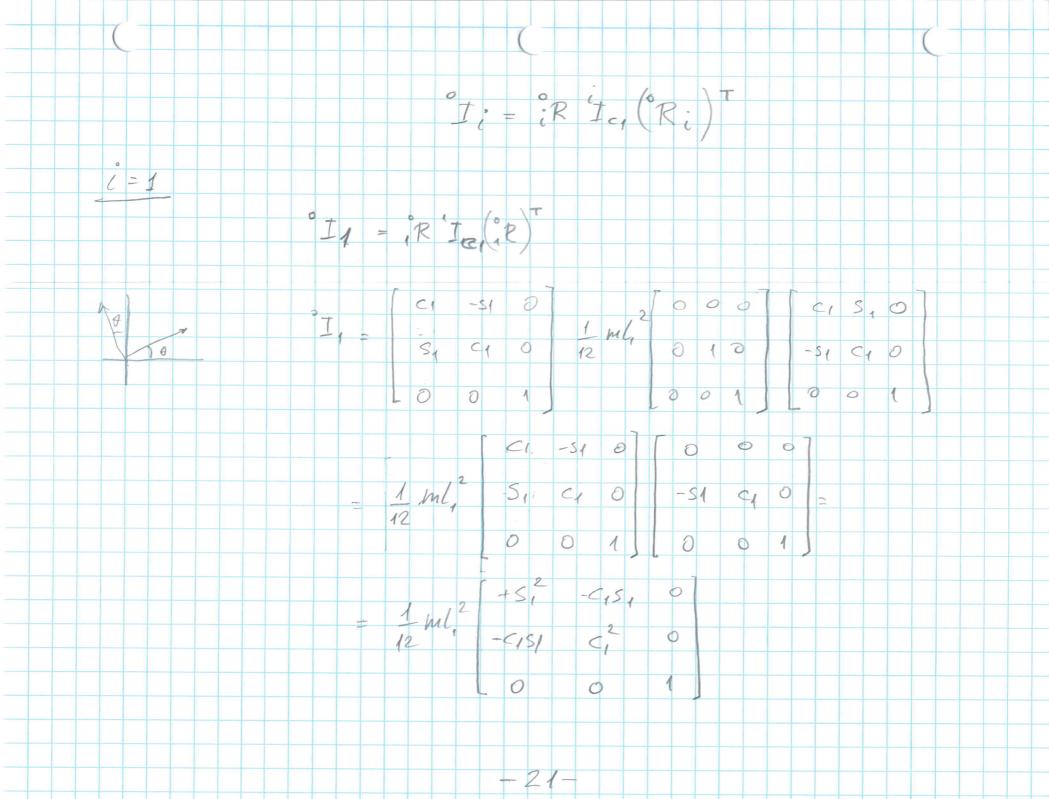
GENERAL FORM OF DYNAMICAL EQUATIONS L = K - U $L = \frac{1}{2} q^T M \dot{q} + \sum_{i=1}^{n} m_i q^T P_{ci}$ Expend the term for the knetic energy into a sum ot scalars Mig - The (i, j) element of the major later inertia materix M $L = \frac{n}{2} \sum_{i=1}^{n} \frac{m_{i}}{j^{2}} \frac{m_{i}}{j^{2}} \frac{q_{i}}{q_{i}} \frac{q_{i}}{q_{j}} \frac{q_{i}}{j^{2}} \frac{q_$ The partial derivative of it with respect to g; Note that the potential energy does not depend on g;

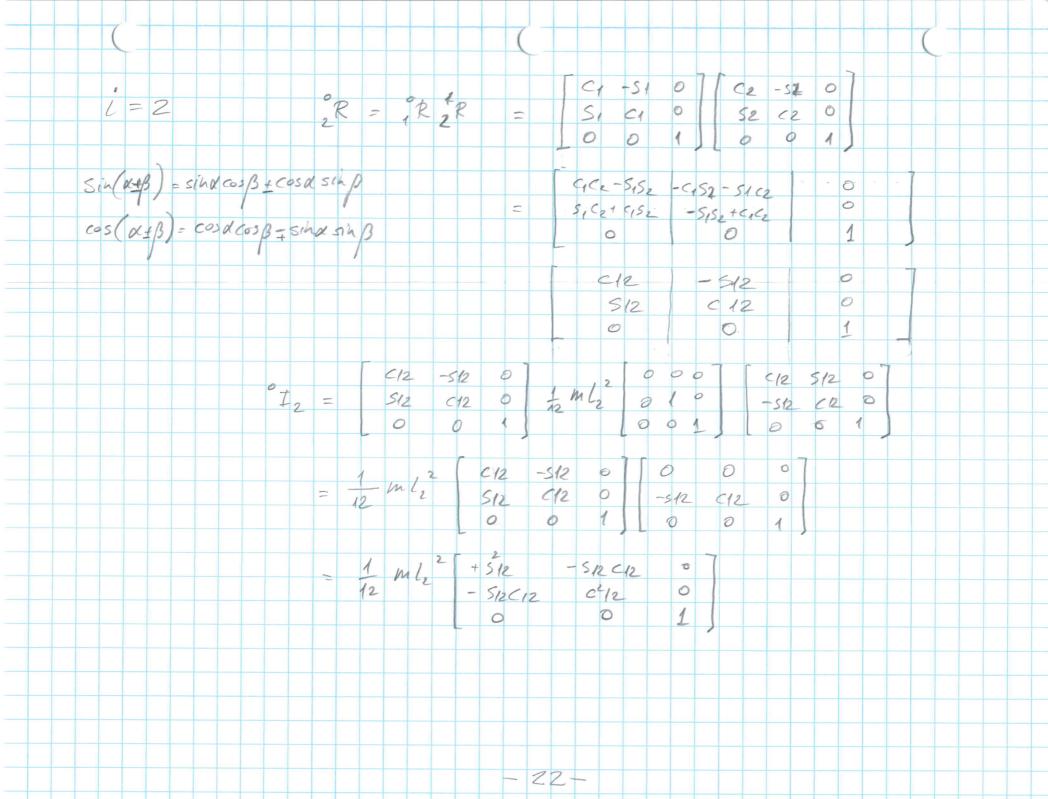
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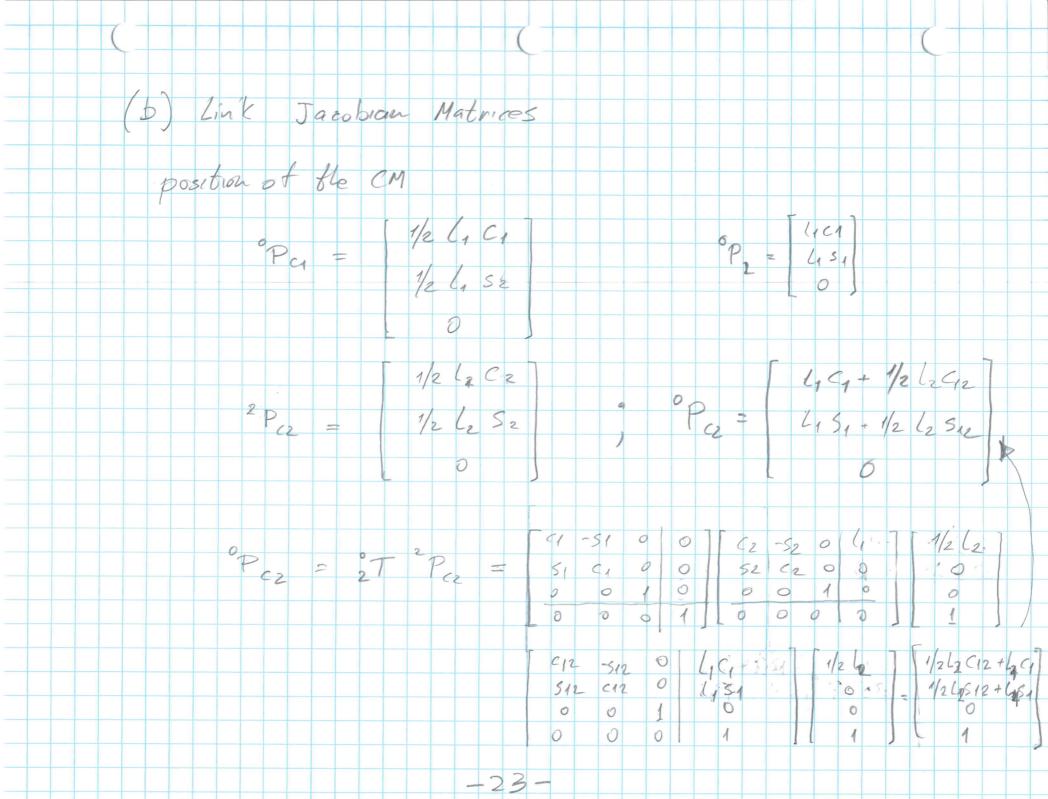


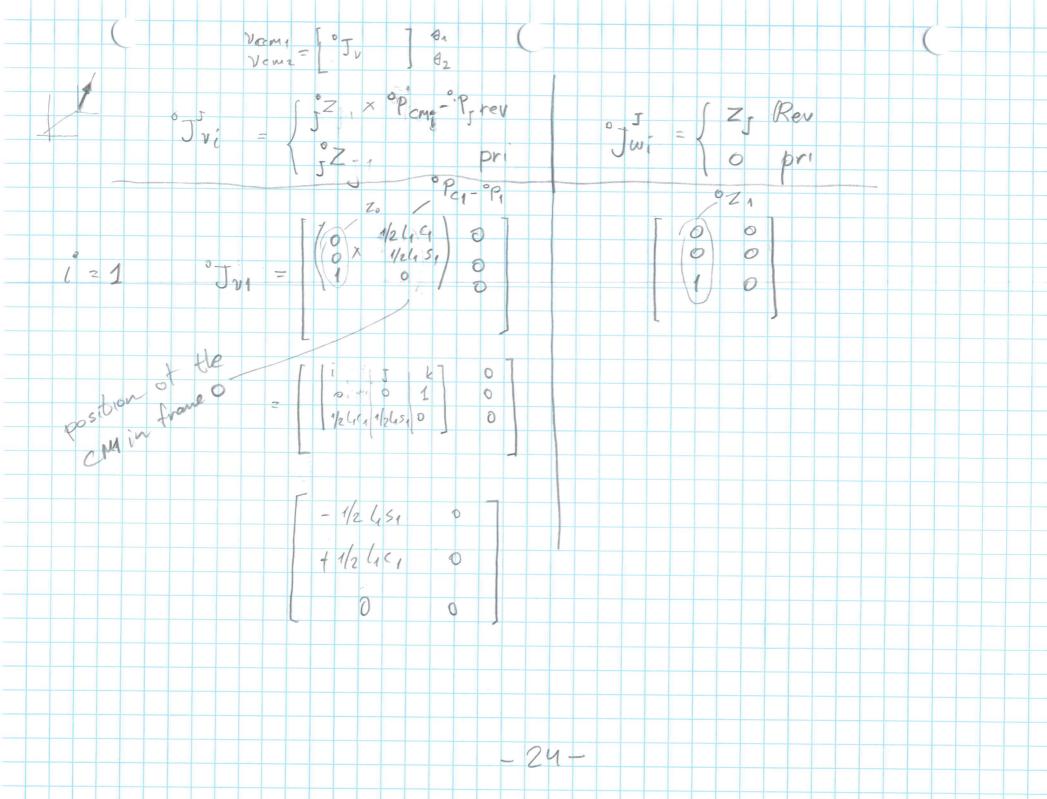
DMIK 95 9K 0L Zhigt 2 Jat N KU + Substitute all equation into the He Eg Lagrange n Nij=1 1. n. n. 93 + Zi Zi 1321 Kel $-1 \frac{2}{2} \frac{\partial M_{FK}}{\partial q_i} \frac{\partial q_k}{\partial f_i} \frac{\partial k}{\partial f_k} - \frac{\partial M_{FK}}{J = 1} \frac{\partial q_j}{\partial f_i} \frac{\partial q_k}{\partial f_i} - \frac{\partial M_{FK}}{J = 1} \frac{\partial q_j}{J = 1} \frac{\partial q_j}{J = 1} \frac{\partial q_j}{J = 1} \frac{\partial q_j}{\partial f_j} \frac{\partial q_k}{\partial f_j} - \frac{\partial M_{FK}}{J = 1} \frac{\partial q_j}{J = 1} \frac{\partial q_j}{\partial f_j} \frac{\partial q_k}{\partial f_j} \frac$ gkgs) Mig OM: T Dqu Mig 9j Signal + Vi + G: Q: n Si j=1 Thertia AV. OMis - 1 OMjk Ogk Z Ogi -E K=1 ĝ, 9k Corio lis contratugal $-\sum_{j=1}^{n} m_j g^j f_{\nu_j}^i$ G Y gravitationa effects 19

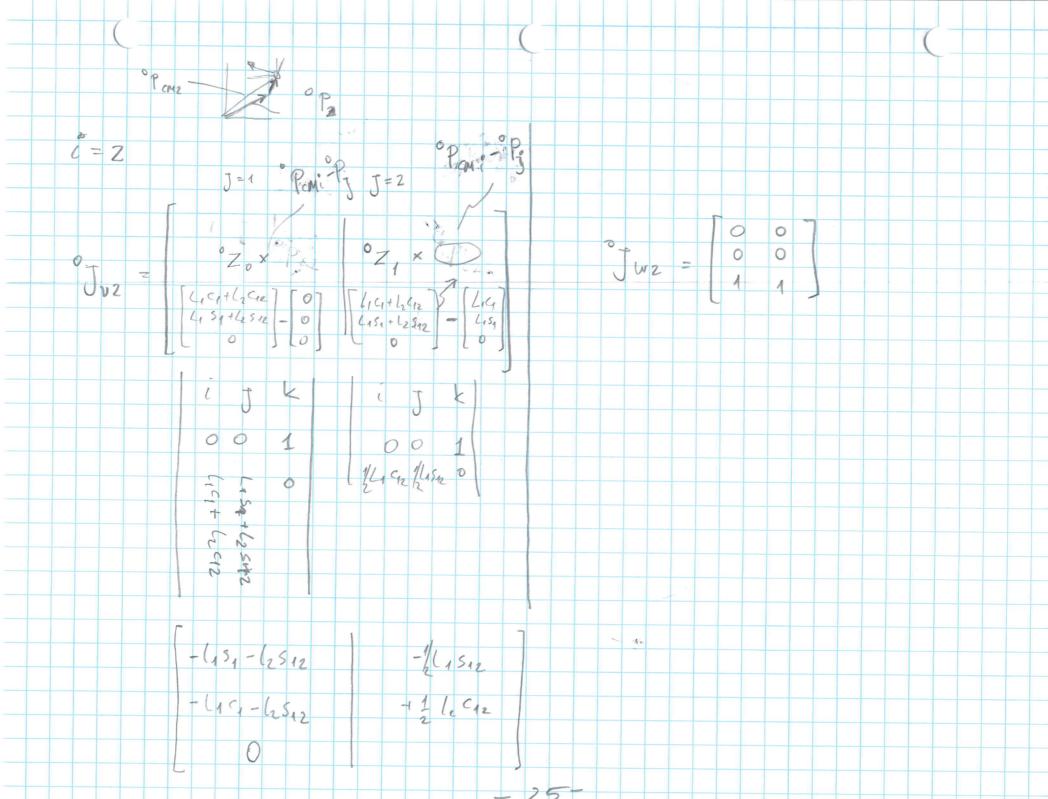


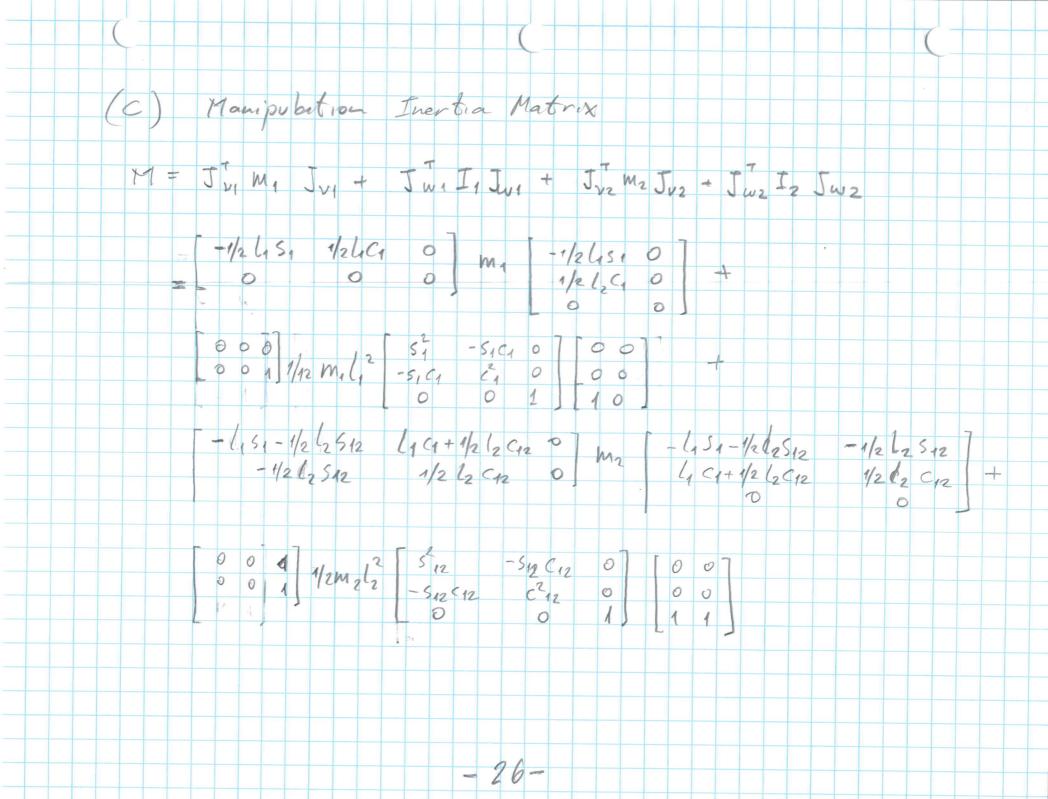


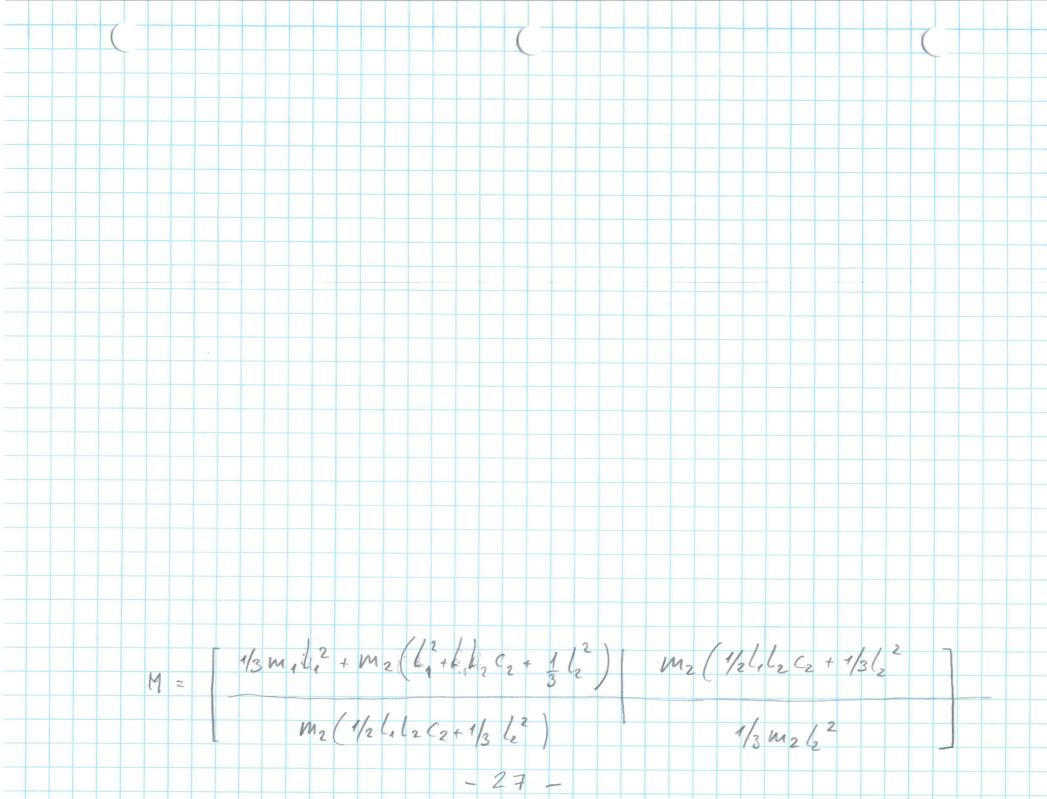


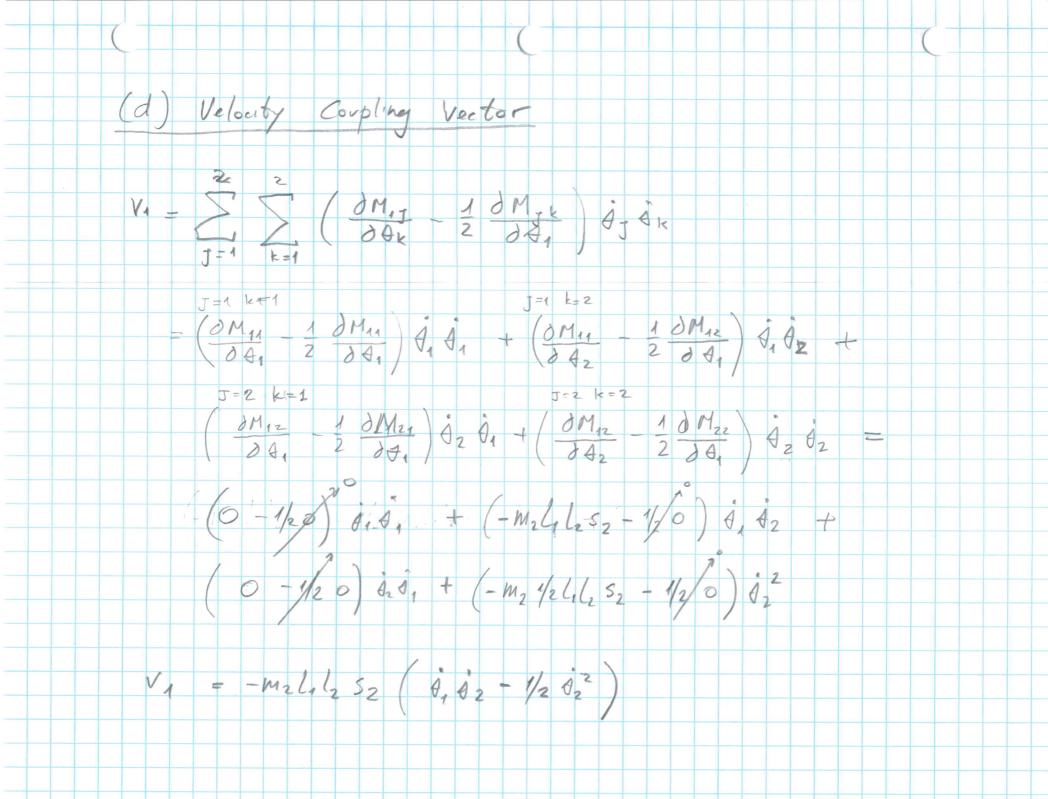


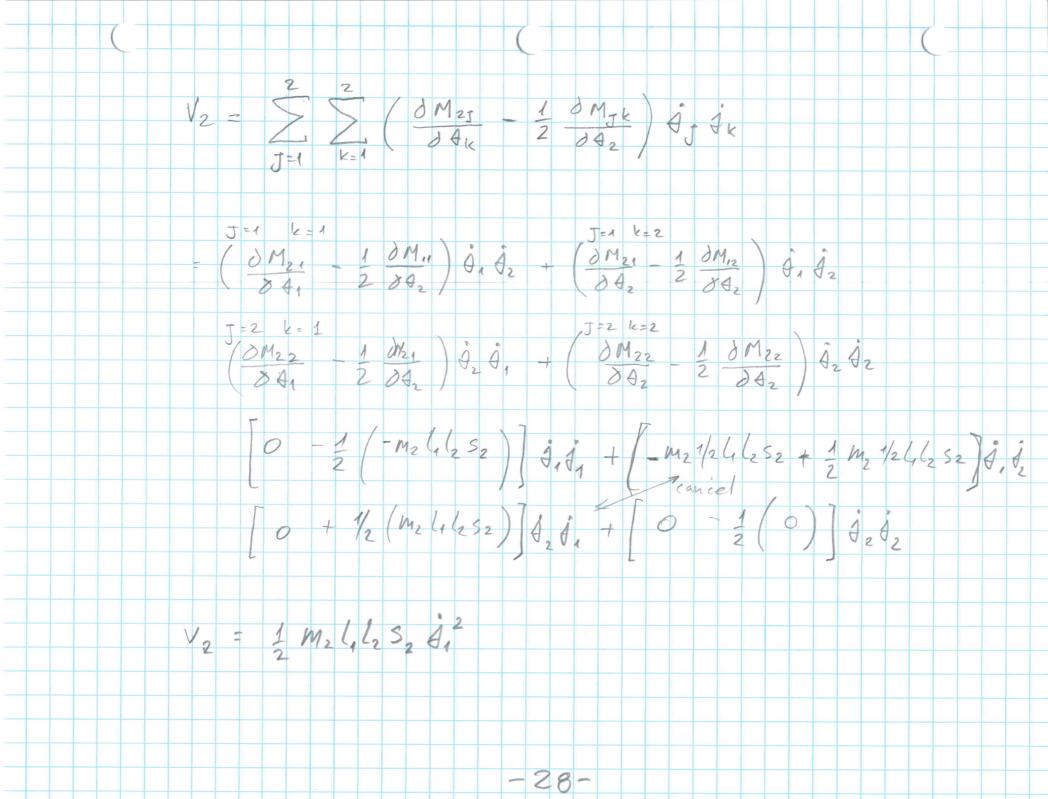


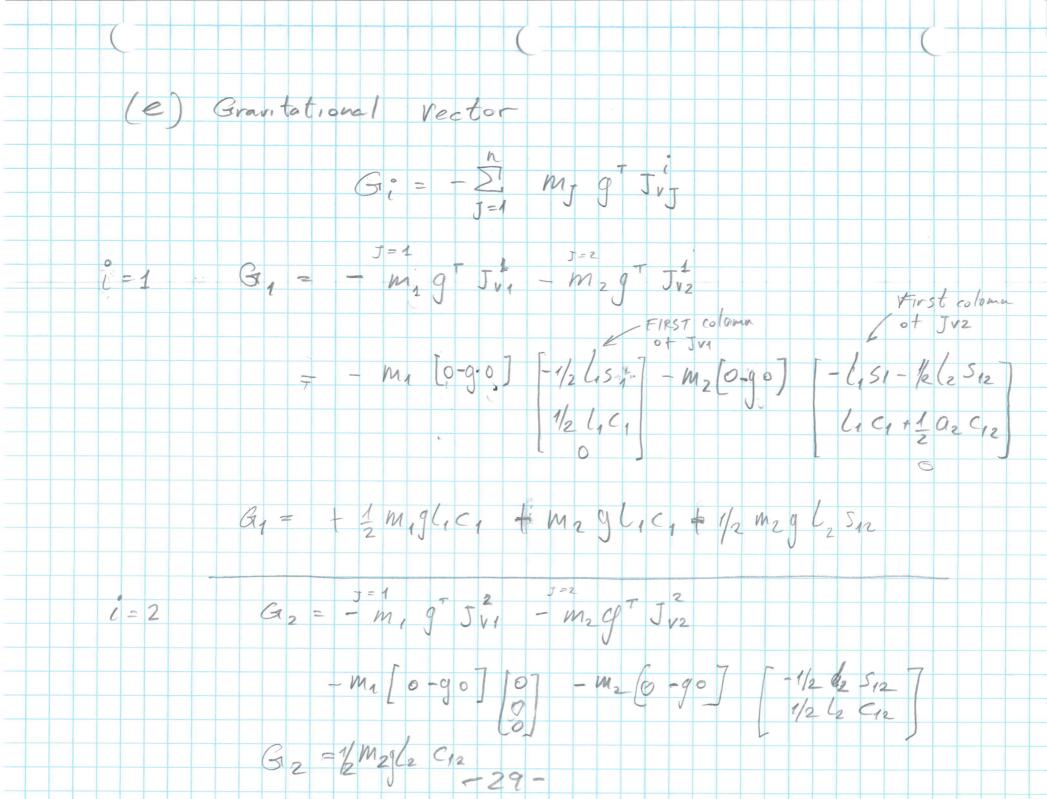


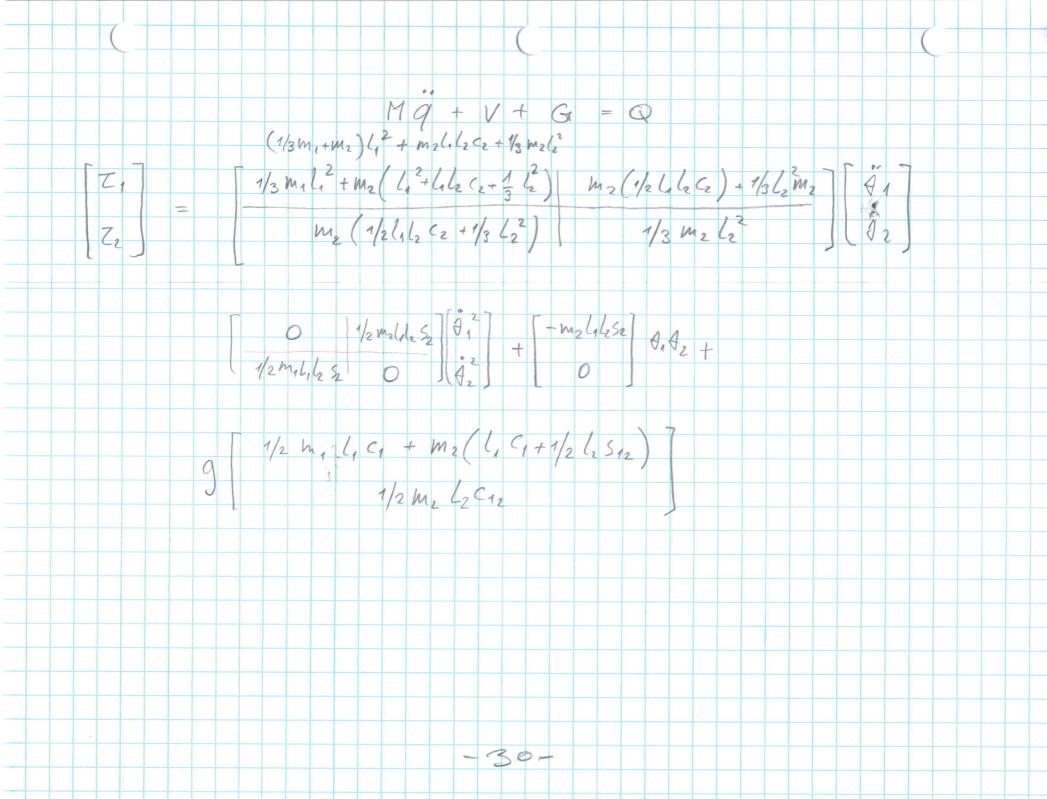


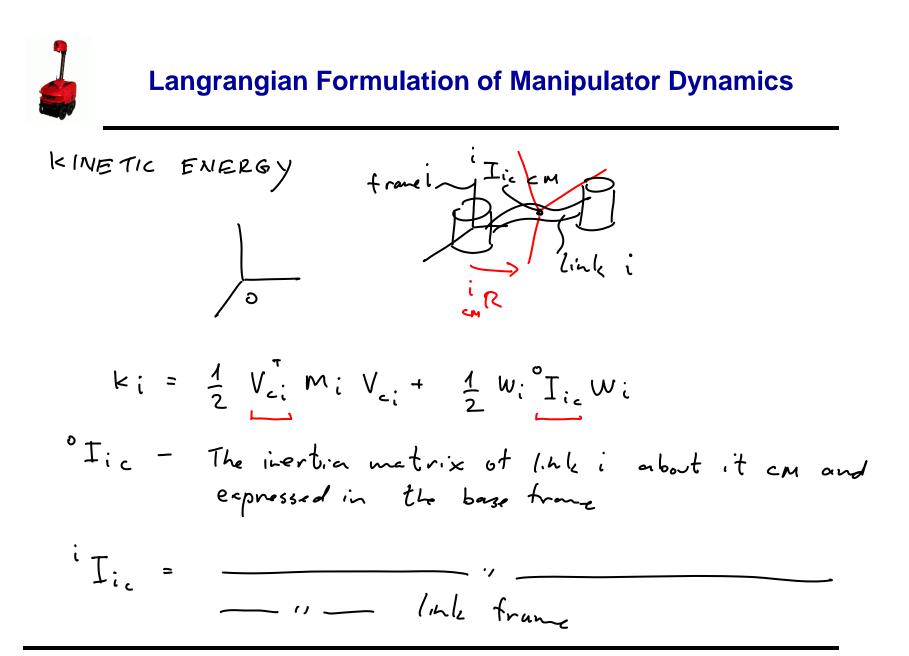




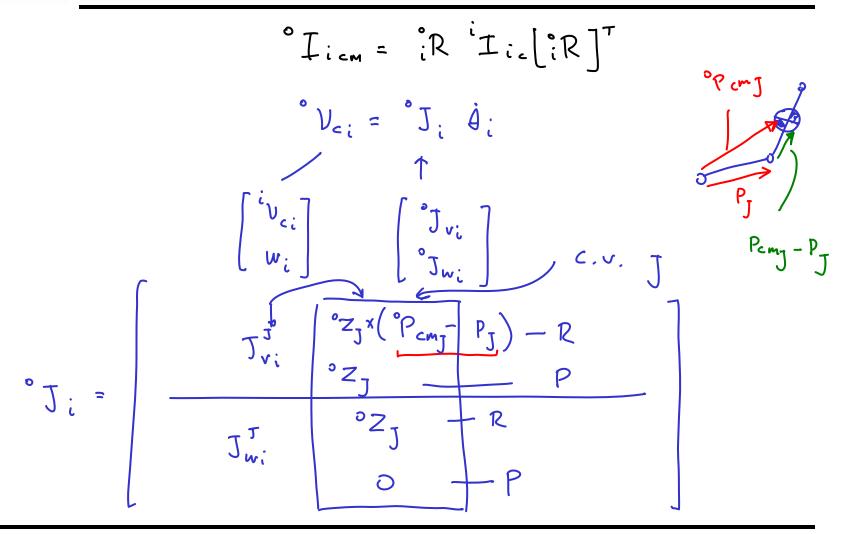




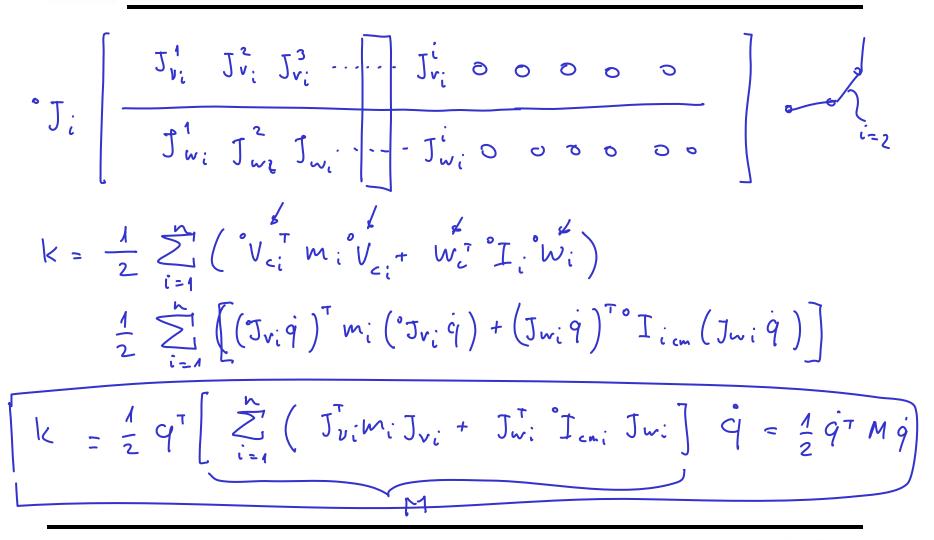




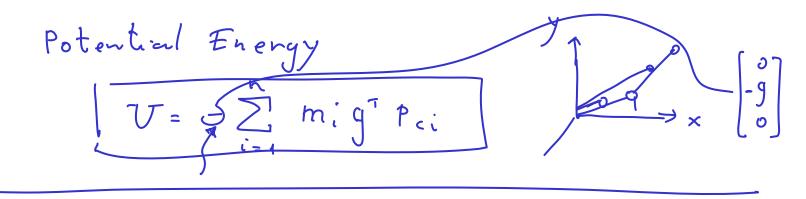




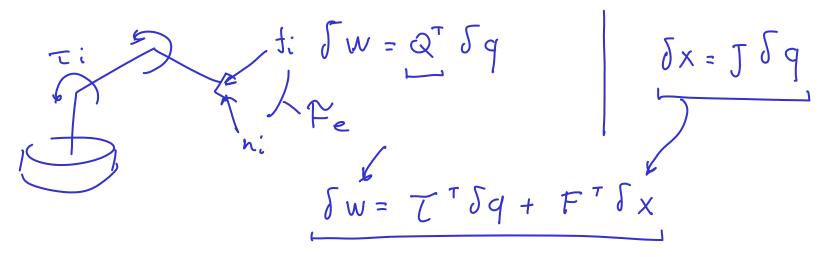




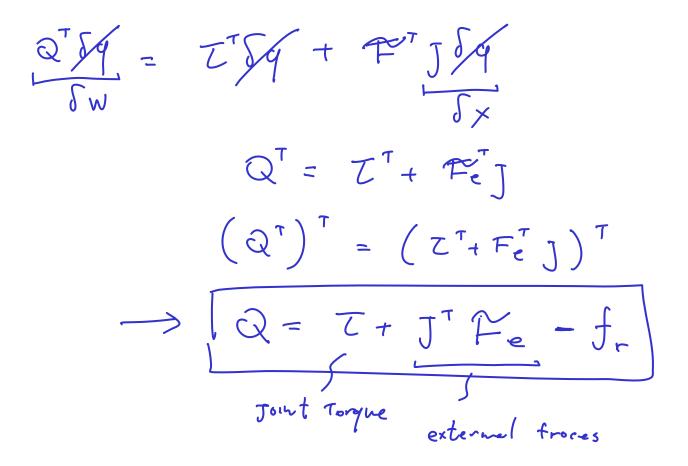




GENERAL FORCES











$$L = k - U$$

$$L = \frac{1}{2} \dot{q}^{T} m \dot{q} \oplus \sum_{i=1}^{n} m_{i} g^{T} P_{ci}$$

$$L = \frac{1}{2} \dot{q}^{T} m \dot{q} \oplus \sum_{i=1}^{n} m_{i} g^{T} P_{ci}$$

$$\longrightarrow \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_{i}} \right) - \frac{\partial L}{\partial q_{i}} = Q_{i}$$

$$\frac{1}{4} \frac{1}{44}$$

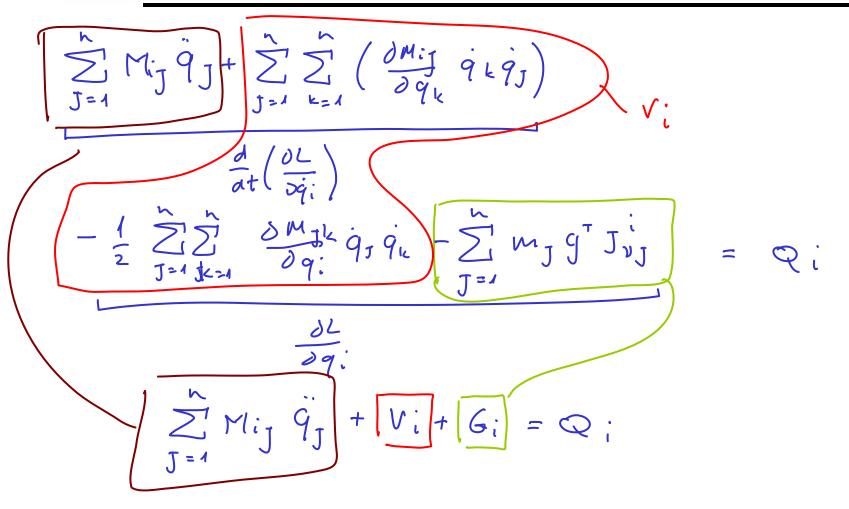




*

$$\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{q}_{i}}\right) = \sum_{j=i}^{n} m_{\dot{t}_{j}} \ddot{q}_{j} + \sum_{j=i}^{n} \left(\frac{d M_{ij}}{dt}\right) \dot{q}_{j} = \frac{d}{dt}\left(\frac{\partial L}{\partial \dot{q}_{i}}\right) = \sum_{j=i}^{n} m_{\dot{t}_{j}} \ddot{q}_{j} + \sum_{j=i}^{n} \left(\frac{d M_{ij}}{dt}\right) \dot{q}_{j} = \frac{d}{dt} + \frac{$$







• It is often convenient to express the dynamic equations of a manipulator in a single equation

$$au = M(heta)\ddot{ heta} + V(heta, \dot{ heta}) + G(heta)$$

where

- $M(\theta)$ Mass matrix (includes inertia terms) *nxn Matrix*
- $V(\hat{\theta}, \dot{\hat{\theta}})$ Centrifugal (square of joint velocity) and Coriolis (product of two different joint velocities) terms *nx1 Vector*
- $G(\theta)$ gravitational terms *nx1 Vector*.



• By rewriting the velocity dependent term $V(\theta, \dot{\theta})$ in a different form, we can write the dynamic equations as

$$\tau = M(\theta)\ddot{\theta} + B(\theta)[\dot{\theta}\dot{\theta}] + C(\theta)[\dot{\theta}^{2}] + G(\theta)$$

where

- $B(\theta)$ Centrifugal coefficients(square of joint velocity) $C(\theta)$ - Coriolis coefficients (product of two different joint velocities)
- This form can be useful for applications using force control. Each of the matrices is a function of manipulator configuration only (that is, joint position) and can be updated at a rate depending on the magnitude of joint changes.



 It can sometimes be desirable to have a relationship between the end effector's Cartesian accelerations and the joint torques. Beginning from the Configuration Space equation

$$\tau = M(\theta)\ddot{\theta} + V(\theta, \dot{\theta}) + G(\theta)$$

• we can substitute the joint moments using our definition of the Jacobian matrix:

$$\tau = J^{T}(\theta)F \qquad F = J^{-T}(\theta)\tau$$
 $\dot{x} = J(\theta)\dot{\theta}$

• By differentiation, we find

$$\ddot{x} = \dot{J}(\theta)\dot{\theta} + J(\theta)\ddot{\theta}$$



Dynamic Equations - Cartesian State Space Equation

• Solving for joint acceleration gives

$$\ddot{\theta} = J^{-1}\ddot{x} - J^{-1}\dot{J}\dot{\theta}$$

• Substitution yields

$$F = J^{-T}\tau = J^{-T}M(\theta)J^{-1}\ddot{x} - J^{-T}M(\theta)J^{-1}\dot{J}\dot{\theta} + J^{-T}V(\theta,\dot{\theta}) + J^{-T}G(\theta)$$

$$F = M_x(\theta)\ddot{x} + V_x(\theta, \dot{\theta}) + G_x(\theta)$$

Where

$$M_{x}(\theta) = J^{-T}M(\theta)J^{-1}$$
$$V_{x}(\theta,\dot{\theta}) = J^{-T}M(\theta)J^{-1}\dot{J}\dot{\theta} + J^{-T}V(\theta,\dot{\theta})$$
$$G_{x}(\theta) = J^{-T}G(\theta)$$

 This equation relates the forces and moments at the end effector to the Cartesian accelerations of the end effector and the manipulator joint positions and velocities.