

4/17/15

MAE182A TA Session #3

P1

Chap3 2nd order D.E.

Goals:

- Homogeneous eqs
 - Superposition
 - Wronskian
 - Complex roots
 - Repeated roots
- Nonhomogeneous Eqs
 - Undetermined Coeff.
 - Variation of parameters
- Mechanical Vibrations
- Forced Vibrations

2nd order D.E.: $\frac{d^2 y}{dt^2} = f(t, y, \frac{dy}{dt})$

linear: $y'' = f(t, y, y') = g(t) - p(t)y' - q(t)y$

$\Rightarrow y'' + p(t)y' + q(t)y = g(t)$ or $p(t)y' + q(t)y + r(t)y = G(t)$

$\left\{ \begin{array}{l} g(t)=0 \Rightarrow \text{homogeneous} \\ g(t) \neq 0 \Rightarrow \text{non homogeneous} \end{array} \right.$

ics: $y(t_0) = y_0$

$y'(t_0) = y'_0$

• Constant coeff^s.

(P2)

$$(1) \Rightarrow ay'' + by' + cy = 0$$

$$\boxed{y = e^{rt}} \Rightarrow (ar^2 e^{rt} + br e^{rt} + ce^{rt}) = 0$$

$$\Rightarrow (ar^2 + br + c) e^{rt} = 0$$

$$\Rightarrow ar^2 + br + c = 0 \quad (\text{char. eq.})$$

$$\Rightarrow r_1, r_2 = ? \quad (2 \text{ roots} \rightarrow 2 \text{ solutions})$$

$$y_1 = e^{r_1 t}, \quad y_2 = e^{r_2 t} \xrightarrow{\text{superposition}} \boxed{y = c_1 y_1 + c_2 y_2}$$

• Linear diff. operator L

$$L(y) = y'' + p(t)y' + q(t)y$$

[Theorem] Existence and Uniqueness (p.146)

$$\text{i.v.p. } y'' + p(t)y' + q(t)y = g(t), \quad y(t_0) = y_0, \quad y'(t_0) = y_0'$$

if p, q, g are continuous on interval I containing t_0 ,
there is exactly one solution existing throughout the interval I

[Theorem] Superposition (p.147)

$$\text{if } y_1 \neq y_2 \text{ are 2 solutions of } L[y] = y'' + p(t)y' + q(t)y = 0^{-(2)}$$

$y = c_1 y_1 + c_2 y_2$ is also a solution for $\forall c_1, c_2$

$$L[c_1 y_1 + c_2 y_2] = c_1 L[y_1] + c_2 L[y_2] \quad (\text{result of linearity of D.E.})$$

• Wronksian

(p3)

$$y = c_1 y_1(t) + c_2 y_2(t) \quad y(t_0) = y_0, \quad y'(t_0) = y'_0$$

$$y(t_0) = c_1 y_1(t_0) + c_2 y_2(t_0) = y_0 \quad \text{--- ①}$$

$$y'(t_0) = c_1 y'_1(t_0) + c_2 y'_2(t_0) = y'_0 \quad \text{--- ②}$$

$$c_1 = \frac{\begin{vmatrix} y_0 & y_2(t_0) \\ y'_0 & y'_2(t_0) \end{vmatrix} = w_1}{\begin{vmatrix} y_1(t_0) & y_2(t_0) \\ y'_1(t_0) & y'_2(t_0) \end{vmatrix} = w}, \quad c_2 = \frac{\begin{vmatrix} y_1(t_0) & y_0 \\ y'_1(t_0) & y'_0 \end{vmatrix} = w_2}{\begin{vmatrix} y_1(t_0) & y_2(t_0) \\ y'_1(t_0) & y'_2(t_0) \end{vmatrix} = w}$$

$$W(t_0) = \begin{vmatrix} y_1(t_0) & y_2(t_0) \\ y'_1(t_0) & y'_2(t_0) \end{vmatrix} = y_1 y'_2 - y'_1 y_2$$

if $W \neq 0$, ① & ② have a unique solution $(c_1, c_2) \Rightarrow y = c_1 y_1 + c_2 y_2$

if $W = 0$, ($w_1 \neq 0$ or $w_2 \neq 0$), no solutions

if $W = 0$, ($w_1 = 0$ & $w_2 = 0$), many solutions

$y = c_1 y_1 + c_2 y_2$ - general solution

y_1, y_2 form a fundamental set of solutions iff $W \neq 0$

• Complex roots of char. eq.

(p4)

$$ay'' + by' + cy = 0$$

$$\boxed{y = e^{rt}} \Rightarrow ar^2 + br + c = 0$$

$$\Rightarrow r = \frac{-b \pm \sqrt{b^2 - 4ac}}{2a}$$

$$\text{If } b^2 - 4ac < 0, \quad r_{1,2} = \lambda \pm i\mu, \quad y_1 = e^{r_1 t}, \quad y_2 = e^{r_2 t}$$

$$\text{Euler's formula: } e^{it} = \cos t + i \sin t$$

$$e^{(\lambda + i\mu)t} = e^{\lambda t} (\cos \mu t + i \sin \mu t)$$

[Theorem 3.2.6] (p.153)

$$L(y) = y'' + p(t)y' + q(t)y = 0, \quad p, q \text{ are real-valued fcn.}$$

If $z = u(t) + i v(t)$ is a complex valued solution,

$u(t)$ & $v(t)$ are also solutions of it

$$y_1 = e^{r_1 t} = e^{(\lambda + i\mu)t} = e^{\lambda t} \cos \mu t + i e^{\lambda t} \sin \mu t$$

[Theorem 3.2.7] Abel's Theorem (p.154)

$$y_1, y_2 \text{ are solutions of } L[y] = y'' + p(t)y' + q(t)y = 0,$$

p, q are continuous on I ,

$$W(y_1, y_2)(t) = c \exp \left[- \int p(t) dt \right]$$

Abel's Theorem

PS

$$y_1'' + p(t)y_1' + q(t)y_1 = 0 \quad \text{--- (1)}$$

$$y_2'' + p(t)y_2' + q(t)y_2 = 0 \quad \text{--- (2)}$$

$$\textcircled{1} \times (-y_2) \Rightarrow -y_2 y_1'' + p(t)y_1'(-y_2) + q(t)y_1(-y_2) = 0$$

$$\textcircled{2} \times y_1 \Rightarrow y_1 y_2'' + p(t)y_1 y_2' + q(t)y_1 y_2 = 0$$

$$(y_1 y_2'' - y_2 y_1'') + (y_1 y_2' - y_1' y_2) p(t) = 0$$

$$W = y_1 y_2' - y_1' y_2$$

$$\Rightarrow W' = y_1 y_2'' - y_1' y_2' - y_1 y_2'' + y_1' y_2'$$

$$\Rightarrow W' + p(t)W = 0$$

$$\Rightarrow \int \frac{dW}{W} = \int -p(t) dt \Rightarrow \ln|W| = -\int p(t) dt \Rightarrow W = c \exp\left[-\int p(t) dt\right]$$

$$y_1 = e^{it} = e^{(\lambda + i\mu)t} = e^{\lambda t} (\cos \mu t + i \sin \mu t) = \underbrace{e^{\lambda t} \cos \mu t}_{u(t)} + i \underbrace{e^{\lambda t} \sin \mu t}_{v(t)}$$

$$\Rightarrow \begin{cases} y_1 = u(t) + i v(t) \\ y_2 = u(t) - i v(t) \end{cases} \Rightarrow u(t) \neq v(t) \text{ are also solutions}$$

$$W(u, v) = uv' - u'v = (e^{\lambda t} \cos \mu t) (\lambda e^{\lambda t} \sin \mu t + \mu e^{\lambda t} \cos \mu t) - (\lambda e^{\lambda t} \cos \mu t - \mu e^{\lambda t} \sin \mu t) e^{\lambda t} \sin \mu t$$

$$\Rightarrow W(u, v) = \mu e^{2\lambda t}$$

$$\mu \neq 0 \Rightarrow W(u, v) \neq 0 \Rightarrow u \neq v \text{ form a fundamental set of solutions}$$

$$\Rightarrow \text{general solution: } y = c_1 u + c_2 v = c_1 e^{\lambda t} \cos \mu t + c_2 e^{\lambda t} \sin \mu t$$

$$\Rightarrow y = e^{\lambda t} (c_1 \cos \mu t + c_2 \sin \mu t)$$

Repeated roots $r_1 = r_2$, $b^2 - 4ac = 0$ (p. 169) #6

$$\Rightarrow r_1 = r_2 = -\frac{b}{2a}, \Rightarrow y_1 = e^{-\frac{b}{2a}t}$$

$$y = e^{r_1 t}, \quad y = v(t) e^{r_1 t} \text{ also a solution}$$

$$y' = v' e^{r_1 t} + v r_1 e^{r_1 t}$$

$$y'' = v'' e^{r_1 t} + \underbrace{v' r_1 e^{r_1 t} + v' r_1 e^{r_1 t}}_{2v' r_1 e^{r_1 t}} + v r_1^2 e^{r_1 t}$$

$$ay'' + by' + cy = 0$$

$$\Rightarrow \left[\underbrace{av''}_{\neq 0} + \underbrace{(2ar_1 + b)v'}_{\neq 0} + \underbrace{(ar_1^2 + br_1 + c)v}_{\neq 0} \right] e^{r_1 t} = 0$$

$$\Rightarrow \begin{cases} 2ar_1 + b = 2a(-\frac{b}{2a}) + b = 0 \\ ar_1^2 + br_1 + c = a(-\frac{b}{2a})^2 + b(-\frac{b}{2a}) + c = \frac{b^2}{4a} - \frac{2b^2}{4a} + \frac{4ac}{4a} = \frac{-(b^2 - 4ac)}{4a} = 0 \end{cases}$$

$$\Rightarrow \underbrace{av''}_{\neq 0} = 0 \Rightarrow v'' = \frac{d}{dt} \left(\frac{dv}{dt} \right) = 0 \Rightarrow \frac{dv}{dt} = \text{const} \Rightarrow v = c_1 + c_2 t$$

$$\therefore y = v e^{r_1 t} = (c_1 + c_2 t) e^{r_1 t} = c_1 \underbrace{e^{r_1 t}}_{\neq 0} + c_2 \underbrace{t e^{r_1 t}}_{\neq 0}$$

• Nonhomogeneous Eqs

(P7)

$$L[y] = y'' + p(t)y' + q(t)y = g(t) \quad (3)$$

[Theorem 3.5.1] p. 176

If $y_1 \neq y_2$ are solutions to the nonhomog. eq. $(y'' + py' + qy = g)$

~~$y_1(t) - y_2(t)$~~ y_1, y_2 are a fund^l set of solutions

$$\text{of } y'' + p(t)y' + q(t)y = 0$$

$$\text{then } y_1 - y_2 = \underbrace{c_1 y_1 + c_2 y_2}_{\text{homog.}}$$

$$\langle p \rangle \quad L[y_1] = g(t), \quad L[y_2] = g(t)$$

$$L[y_1] - L[y_2] = g(t) - g(t) = 0$$

$$\Rightarrow L[y_1 - y_2] = 0.$$

[Theorem 3.5.2]

General solution to nonhomog. eq.

$$y(t) = \phi(t) = \underbrace{c_1 y_1 + c_2 y_2}_{\text{homog. solut.}} + \underbrace{Y_p(t)}_{\text{particular solution}}$$

How to find $Y_p(t)$?

Undetermined coeff.

- choose solution of the same form as nonhomog. term $g(t)$

- polynomial
- exponential
- cos, sin

- if Y_p has the same form as homog. solution.

multiply by t until there is no duplication

$L[y]$	$[Y_p]$
$5t^2$	$At^2 + Bt + C$
$\cos 3t$	$A \cos 3t + B \sin 3t$
$8e^{5t}$	Ae^{5t}
$t^2 \sin t$	$(At^2 + Bt + C)(\cos t + \sin t)$
$t + e^{5t}$	$(At + B) + Ce^{5t}$

ex p. 17, p. 184

(P9)

$$y'' - 2y' + y = te^t + 4.$$

1° Homog. sol: $y'' - 2y' + y = 0$

$$y = e^{rt} \Rightarrow r^2 - 2r + 1 = 0 \Rightarrow r_{1,2} = 1$$

$$y_h = c_1 e^t + c_2 t e^t$$

2° $g_1(t) = te^t$

Guess $Y_{p1} = (At+B)e^t$, but te^t is a homog. sol.

$$\xrightarrow{*t^2} Y_{p1} = (At^3 + Bt^2)e^t$$

$$Y'_{p1} = (3At^2 + 2Bt)e^t + (At^3 + Bt^2)e^t$$

$$Y''_{p1} = (6At + 2B)e^t + (3At^2 + 2Bt)e^t + (3At^2 + 2Bt)e^t + (At^3 + Bt^2)e^t$$

$$\Rightarrow Y_{p1} = \frac{t^3 e^t}{6} \quad (A = \frac{1}{6}, B = 0)$$

3° $g_2(t) = 4$

$$Y_{p2} = c \Rightarrow 0 - 0 + c = 4 \Rightarrow c = 4$$

$$\therefore y = y_h + y_{p1} + y_{p2} = c_1 e^t + c_2 t e^t + \frac{t^3 e^t}{6} + 4.$$

Variation of Parameters

(P. 0)

$$y'' + p(t)y' + q(t)y = g(t) \quad [4]$$

const. coeff. $\Rightarrow ay'' + by' + cy = 0$

general sol. $\Rightarrow y = c_1 y_1(t) + c_2 y_2(t)$

* $p(t), q(t)$ are not const.

[Note] $p(t) \neq q(t)$ are not const but the fns of t

[Theorem 3.6.1]
(p. 189) $y = u_1(t) y_1(t) + u_2(t) y_2(t)$

$$\Rightarrow y'', y'$$

$$\Rightarrow u_1(t) = - \int \frac{y_2 g(t)}{W(y_1, y_2)} dt + c_1$$

$$u_2(t) = \int \frac{y_1 g(t)}{W(y_1, y_2)} dt + c_2$$

a particular solution for [4]

$$Y(t) = -y_1(t) \int_{t_0}^t \frac{y_2 g(s)}{W(y_1, y_2)} ds + y_2(t) \int_{t_0}^t \frac{y_1(s) g(s)}{W(y_1, y_2)} ds$$

General solution:

$$y = c_1 y_1(t) + c_2 y_2(t) + Y(t)$$

ex $5y'' + (t)y' + 7 = t^2 \Rightarrow g(t) = \frac{t^2}{5}$

$$y'' + \left(\frac{t}{5}\right)y' + \frac{7}{5} = \frac{t^2}{5} = g(t)$$

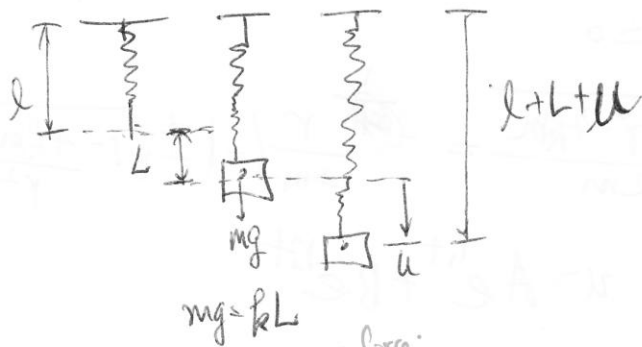
[Note]

$$g(t) = \frac{t^2}{5}$$

$$g(t) \neq t^2!$$

Mechanical Vibrations

P11



FBD $\vec{F} = -k(l+u)$ (spring force)
 $\vec{F}_d = -ru'$ (damping force)
 $\vec{F}_g = mg$ (gravity)

$$\sum \vec{F} = \vec{F}_g + \vec{F}_s + \vec{F}_d + \vec{F}$$

$$\Rightarrow m\ddot{u} = mg - k(l+u) - ru' + F(t)$$

$$\Rightarrow m\ddot{u} = -ku - ru' + F(t)$$

$$\Rightarrow m\ddot{u} + ru' + ku = F(t)$$

Free vibration. $F(t) = 0 \Rightarrow m\ddot{u} + ru' + ku = 0$

[Case 1] Undamped vibrations $F(t) = 0, r = 0$

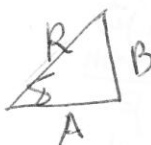
$$m\ddot{u} + ku = 0$$

$$u = e^{rt} \Rightarrow mr^2 + k = 0 \Rightarrow r = \pm i\sqrt{\frac{k}{m}} = \pm i\omega_0$$

$$\Rightarrow u = A \cos \omega_0 t + B \sin \omega_0 t \quad \text{where } \omega_0^2 = \frac{k}{m}$$

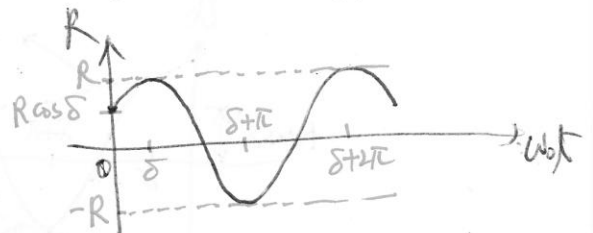
$$\begin{cases} A = R \cos \delta \\ B = R \sin \delta \end{cases}$$

$$\Rightarrow u = R \cos \omega_0 t \cos \delta + R \sin \omega_0 t \sin \delta = R \cos(\omega_0 t - \delta)$$



$$\Rightarrow R = \sqrt{A^2 + B^2}, \quad \tan \delta = \frac{B}{A}$$

↑ ↑
amplitude phase



$$\text{Period} = T = \frac{2\pi}{\omega_0} = 2\pi \left(\frac{m}{k} \right)^{\frac{1}{2}}, \quad \omega_0 = \text{natural freq.}$$

✓ [case 2] Damped $r \neq 0$

(P12)

$$mu'' + ru' + ku = 0$$

$$r_1, r_2 = \frac{-r \pm \sqrt{r^2 - 4km}}{2m} = \frac{r}{2m} (-1 \pm \sqrt{1 - \frac{4km}{r^2}})$$

$$\begin{cases} r^2 - 4km > 0 \Rightarrow u = Ae^{r_1 t} + Be^{r_2 t} \\ r^2 - 4km = 0 \Rightarrow u = (A + Bt) e^{-\frac{rt}{2m}} \\ r^2 - 4km < 0 \Rightarrow u = e^{-\frac{rt}{2m}} (A \cos \mu t + B \sin \mu t) \end{cases}$$

$$\mu = \frac{(4km - r^2)^{\frac{1}{2}}}{2m} > 0$$

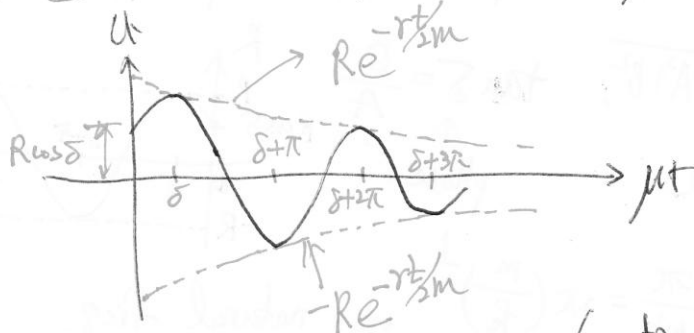
$$r^2 - 4km \geq 0, \quad \frac{r^2 - 4km}{r^2} < \frac{r^2}{r^2} = 1 \Rightarrow \sqrt{\frac{r^2 - 4km}{r^2}} < 1 \Rightarrow r_1, r_2 < 0$$

$$r^2 - 4km < 0 \quad -\frac{r}{2m} < 0 \Rightarrow \text{real part} < 0$$

all 3 solutions: $t \rightarrow \infty, \quad u \rightarrow 0$

$$r^2 - 4km < 0 \Rightarrow u = e^{-\frac{rt}{2m}} (A \cos \mu t + B \sin \mu t)$$

$$\begin{cases} A = R \cos \delta \\ B = R \sin \delta \end{cases} \Rightarrow u = R e^{-\frac{rt}{2m}} \cos(\mu t - \delta)$$



quasi freq μ

$$\frac{\mu}{\omega_0} = \frac{(4km - r^2)^{\frac{1}{2}}}{(km)^{\frac{1}{2}}}$$

$$= (1 - \frac{r^2}{4km})^{\frac{1}{2}} \approx 1 - \frac{r^2}{8km}$$

(not a period but oscillating at μ)

• Forced Vibration (§3.8)

P13

$$m u''(t) + \gamma u' + k u = F(t) = F_0 \cos \omega t$$

$$u = c_1 u_1(t) + c_2 u_2(t) + A \cos \omega t + B \sin \omega t$$

$$= \underbrace{u_h(t)}_{\text{transient}} + \underbrace{u_p(t)}_{\text{steady}}$$

- transient homogen. solution dies out but it allows us to satisfy i.c.s
- steady solution.

$$u(t) = R \cos(\omega t - \delta)$$

$$R = \frac{F_0}{\Delta}, \quad \cos \delta = \frac{m(\omega_0^2 - \omega^2)}{\Delta}, \quad \sin \delta = \frac{\gamma \omega}{\Delta}$$

$$\Delta = \sqrt{m^2(\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2}, \quad \omega_0^2 = \frac{k}{m}$$

$$\frac{Rk}{F_0} = \frac{1}{\left[\left(1 - \frac{\omega^2}{\omega_0^2}\right)^2 + \gamma^2 \frac{\omega^2}{\omega_0^4} \right]^{1/2}}, \quad T = \frac{\gamma^2}{mk}$$

low freq: $\omega \rightarrow 0, \quad R = \frac{F_0}{k}$

high freq: $\omega \rightarrow \infty, \quad R \rightarrow 0$

Max amplitude, $\omega = \omega_{\max} = ?$

$$\frac{d}{d\omega} \left[\frac{Rk}{F_0} \right] = 0$$

P4

$$\omega_{\max}^2 = \omega_0^2 - \frac{r^2}{2m^2} = \omega_0^2 \left(1 - \frac{r^2}{2mk} \right)$$

[Note] 1. $\omega_{\max} < \omega_0$

2. $r \ll 1 \quad \therefore \omega_{\max} \approx \omega_0$

3. $R_{\max} = \frac{F_0}{r\omega_0 \sqrt{1 - \frac{r^2}{4mk}}} \approx \frac{F_0}{r\omega_0} \left(1 + \frac{r^2}{8mk} \right)$

4. ~~$\frac{\omega}{\omega_0} \approx 1$~~ $r \ll 1 \quad R_{\max} \approx \frac{F_0}{r\omega_0}$

~~F_0 is small, but $R_{\max}$ is quite large~~

r is small $\Rightarrow R_{\max} \uparrow$ (Resonance)

