

Chapter 3 Second Order DE

$$\frac{d^2 y}{dt^2} = f\left(t, y, \frac{dy}{dt}\right) = \underbrace{g(t) - p(t) \frac{dy}{dt} - q(t) y}_{\text{if linear}}$$

$$y(t_0) = y_0$$

$$y'(t_0) = y'_0$$

$$y'' + p(t)y' + q(t)y = g(t)$$

↑
non homogeneous term

Const. coefficients

$$ay'' + by' + cy = 0 \rightarrow \text{seek sol. of form } y = e^{rt}$$

$$\underbrace{(ar^2 + br + c)}_{\substack{\text{must equal} \\ \text{zero}}} e^{rt} = 0 \quad \text{never zero}$$

$$\frac{ar^2 + br + c = 0}{\text{characteristic eq.}}$$

Second order eq. gives 2 roots \rightarrow 2 solutions

$$y_1 = e^{r_1 t}, \quad y_2 = e^{r_2 t}$$

$$\rightarrow \boxed{y = C_1 y_1 + C_2 y_2}$$

general solution

Linear operator, L

$$L[y] = y'' + p(t)y' + q(t)y$$

Theorems

Existence + Uniqueness - pg 146

$$y'' + p(t)y' + q(t)y = g(t) \quad y(t_0) = y_0 \quad y'(t_0) = y_0'$$

if p, q, g is continuous on interval I containing t_0
There is exactly one solution

Superposition pg 147

If y_1, y_2 are two sol. of

$$L[y] = y'' + p(t)y' + q(t)y = 0$$

$y = C_1 y_1 + C_2 y_2$ is also a solution

$$L[C_1 y_1 + C_2 y_2] = C_1 L[y_1] + C_2 L[y_2] \quad \text{pg 148}$$

Result of linearity of DE!

Wronskian

for two solutions y_1, y_2

we know $y = C_1 y_1(t) + C_2 y_2(t)$ is solution.

We can always find C_1, C_2 if the
Wronskian Determinant of y_1, y_2 is not zero
at t_0

$$W = y_1 y_2' - y_1' y_2 = \det \begin{vmatrix} y_1 & y_2 \\ y_1' & y_2' \end{vmatrix}$$

If $W(t_0) \neq 0$ then $y = C_1 y_1 + C_2 y_2$
contains every solution.

$$y = C_1 y_1 + C_2 y_2 \quad \text{— general solution}$$

$y_1 + y_2$ form fundamental set
of solutions

Complex Roots

$$r_{1,2} = \frac{-b \pm \sqrt{b^2 - 4ac}}{2a}$$

$$b^2 - 4ac < 0$$

$$r_{1,2} = \alpha \pm i\mu \quad y_1 = e^{r_1 t} \quad y_2 = e^{r_2 t}$$

We can use Taylor Series to show

$$e^{it} = \cos t + i \sin t$$

$$e^{(\alpha + i\mu)t} = e^{\alpha t} e^{i\mu t} = e^{\alpha t} (\cos \mu t + i \sin \mu t)$$

Since IVP has only real coefficient we use Theorem 3.2.6 which states the real + imaginary parts of a sol. are also solutions

$$y_1 = e^{\alpha t} \cos \mu t + i \sin \mu t$$

$$y_2 = e^{\alpha t} \cos \mu t - i \sin \mu t$$

$$\rightarrow y = e^{\alpha t} (c_1 \cos \mu t + c_2 \sin \mu t)$$

Repeated Roots $r_1 = r_2$

$$y = c_1 e^{rt} + c_2 t e^{rt}$$

Nonhomogeneous Equations

$$L[y] = y'' + p(t)y' + q(t)y = g(t)$$

Theorem 3.5.1

If y_1, y_2 sol. to non homogeneous eq.

$$y_1 - y_2 = \underbrace{c_1 y_1 + c_2 y_2}_{\text{homogeneous solution}}$$

General solution to non homogeneous eq.

$$y = \phi(t) = \underbrace{c_1 y_1(t) + c_2 y_2(t)}_{\text{homogeneous}} + \underbrace{y_p(t)}_{\substack{\uparrow \text{particular} \\ \text{solution}}}$$

To find Particular sol.

Method of Undetermined Coefficients.

- choose solution of the same form as non homogeneous term.

- polynomial
- exponential.
- cos, sin.

- If y_p same form as homogeneous solution multiply by t . until there is no duplication

$\mathcal{L}[y]$	$[Y_P]$
s^2	$At^2 + Bt + C$
$\cos 3t$	$A \cos 3t + B \sin 3t$
$8e^{5t}$	Ae^{5t}
$t^2 \sin 2t$	$(At^2 + Bt + C)(\cos 2t + \sin 2t)$
$t + e^{5t}$	$(At + B) + (e^{5t})$

↑
 For sum you can solve as two separate problems!

prob 17 pg 184

$$y'' - 2y' + y = te^t + 4$$

Homogeneous equation $\rightarrow r^2 - 2r + 1 = 0$
 $r_1 = r_2 = 1$

$$y_h(t) = c_1 e^t + c_2 t e^t$$

Consider $g_1(t) = te^t$, we would guess $(At + b)e^t$
 but te^t is a homo solution

- set $y_1 = At^2 e^t + Bt^3 e^t$ (multiplied by t^2)

- Plug into non homogeneous eq. and solve for coefficients

$$[At^2 e^t + Bt^3 e^t]'' - 2[At^2 e^t + Bt^3 e^t]' + At^2 e^t + Bt^3 e^t = te^t$$

collect like terms and set coeff = 0

$$y_1 = \frac{t^3 e^t}{6}$$

for $g_2(t) = 4$

- set $y_2 = C$

- plug in

$$0 - 0 + C = 4$$

$$C = 4$$

$$y_2 = 4$$

$$y = C_1 e^t + C_2 t e^t + \frac{t^3 e^t}{6} + 4$$

Variation of Parameters

- general method

- $g(t)$ can have any form

$$y(t) = -y_1(t) \int_{t_0}^t \frac{y_2(s)g(s)}{W(y_1, y_2)(s)} ds + y_2(t) \int_{t_0}^t \frac{y_1(s)g(s)}{W(y_1, y_2)(s)} ds$$

$$y = C_1 y_1(t) + C_2 y_2(t) + y(t)$$

- don't forget bounds

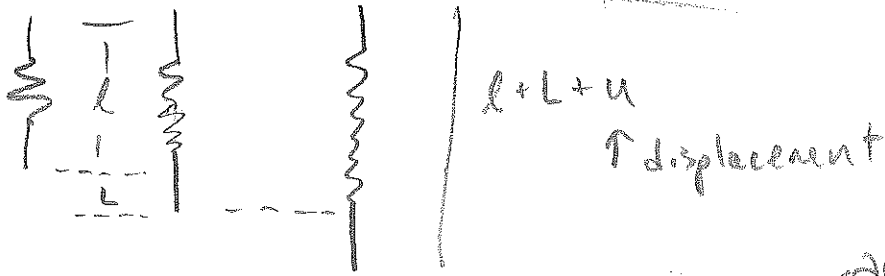
- Integrals may be difficult to evaluate.

- equation must be in correct form to use above eq.

$$5y'' + ty' + 7 = t^2 \Rightarrow g(t) = \frac{t^2}{5}$$

$g(t) \neq t^2$!

Mechanical Vibrations



$$m u''(t) + \gamma u'(t) + k u(t) = F(t)$$

← derived in Discussion 1

↑ Forcing function

Free vibration $F(t) = 0$

Undamped. $\gamma = 0 \rightarrow m u'' + k u = 0$

$$m r^2 + k = 0 \rightarrow r = \pm i \sqrt{\frac{k}{m}}$$

$$u = A \cos \omega_0 t + B \sin \omega_0 t \quad \text{where } \omega_0^2 = \frac{k}{m}$$

Lets write in different form

$$u = R \cos(\omega_0 t - \delta)$$

$$= R \cos \delta \cos \omega_0 t + R \sin \delta \sin \omega_0 t$$

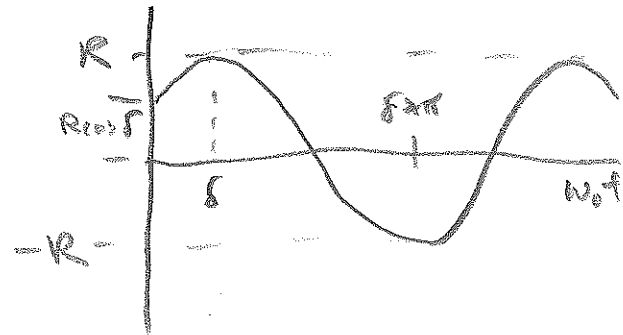
$$A = R \cos \delta$$

$$B = R \sin \delta$$

amplitude ↑ phase

$$R = \sqrt{A^2 + B^2}$$

$$\tan \delta = \frac{B}{A}$$



Period $T = \frac{2\pi}{\omega_0} = 2\pi \left(\frac{m}{k}\right)^{1/2}$

↑ natural frequency

Damped

$$m\ddot{u} + \gamma\dot{u} + ku = 0$$

$$r_1, r_2 = \frac{-\gamma \pm \sqrt{\gamma^2 - 4km}}{2m} = \frac{\gamma}{2m} \left(-1 \pm \sqrt{1 - \frac{4km}{\gamma^2}} \right)$$

$$\left[\begin{array}{l} \gamma^2 - 4km > 0 \end{array} \right. \rightarrow u = Ae^{r_1 t} + Be^{r_2 t} \quad \text{roots are negative}$$

$$\left[\begin{array}{l} \gamma^2 - 4km = 0 \end{array} \right. \quad u = (A + Bt) e^{-\gamma t / 2m}$$

$$\left[\begin{array}{l} \gamma^2 - 4km < 0 \end{array} \right. \quad u = e^{-\gamma t / 2m} (A \cos \mu t + B \sin \mu t) \quad \text{negative real part}$$

$$\mu = \frac{(4km - \gamma^2)^{1/2}}{2m} > 0$$

Negative exponentials

$$u \rightarrow 0, \quad t \rightarrow \infty$$

note! γ value doesn't matter

γ^2 relative to $4km$ matters

quasi frequency - μ

- motion not periodic but oscillates at μ .

$$\frac{\mu}{\omega_0} = \left(1 - \frac{\gamma^2}{4km} \right)^{1/2}$$

Damping reduces frequency!

Forced Vibrations

$$m u''(t) + \gamma u' + k u = F(t) = F_0 \cos \omega t$$

$$u = C_1 u_1(t) + C_2 u_2(t) + A \cos \omega t + B \sin \omega t$$
$$= u_c(t) + U(t)$$

↑ ↑
transient steady

~~*~~ We know transient homogeneous solution dies out but it allows us to satisfy initial conditions - convenient to express $U(t)$ as single term

$$U(t) = R \cos(\omega t - \delta)$$

$$R = \frac{F_0}{\Delta} \quad \cos \delta = \frac{m(\omega_0^2 - \omega^2)}{\Delta} \quad \sin \delta = \frac{\gamma \omega}{\Delta}$$

$$\Delta = \sqrt{m^2(\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2} \quad \omega_0^2 = \frac{k}{m} \quad \Gamma = \frac{\gamma^2}{m k}$$

gives

$$\frac{R k}{F_0} = \frac{1}{\left[\left(1 - \frac{\omega^2}{\omega_0^2} \right)^2 + \Gamma \frac{\omega^2}{\omega_0^2} \right]^{1/2}}$$

$\omega \rightarrow 0$ $R = \frac{F_0}{k}$ low frequency

$\omega \rightarrow \infty$ $R \rightarrow 0$ high frequency

max m between?

differentiate and set = 0

$$W_{\max} = W_0^2 \left(1 - \frac{\delta^2}{2mk} \right)$$

$$W_{\max} < 0$$

$$R_{\max} = \frac{F_0}{\delta W_0 \sqrt{1 - \frac{\delta^2}{4mk}}}$$

Resonance
Frequency